

Integration of Magnons into Superconducting Quantum Circuits

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European Research Council
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Quantum
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INSTITUTO DE NANOCIENCIA
Y MATERIALES DE ARAGÓN

❖ Light-matter interaction and circuit QED

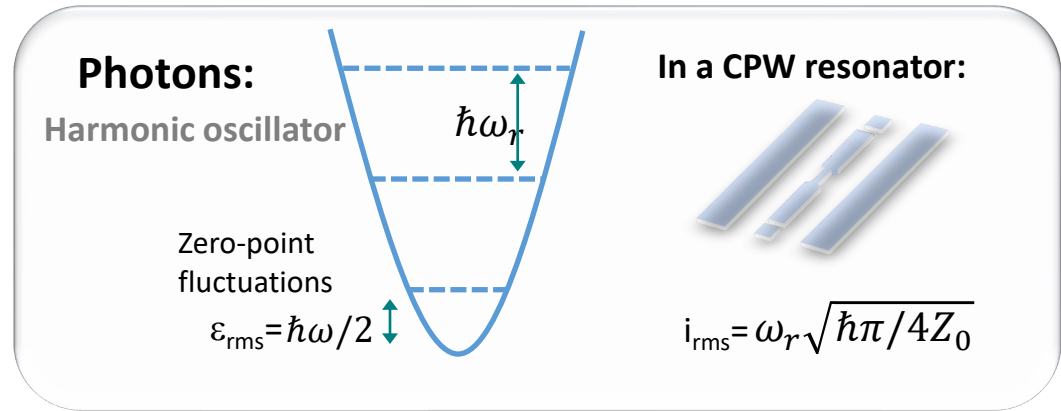
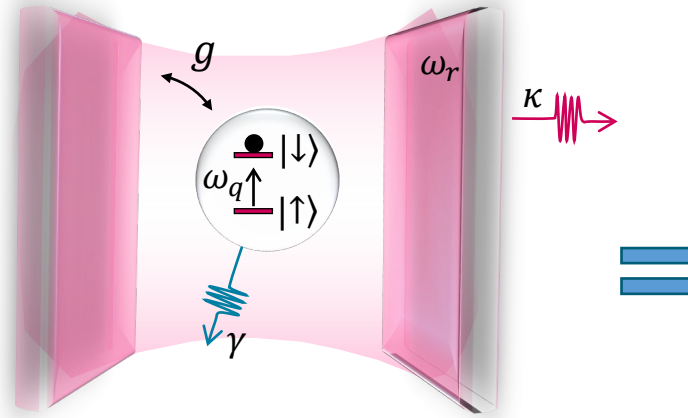
❖ Nanomagnets + **superconducting resonators**

- 😊 Coupling to spins
- 😊 Coupling to magnons

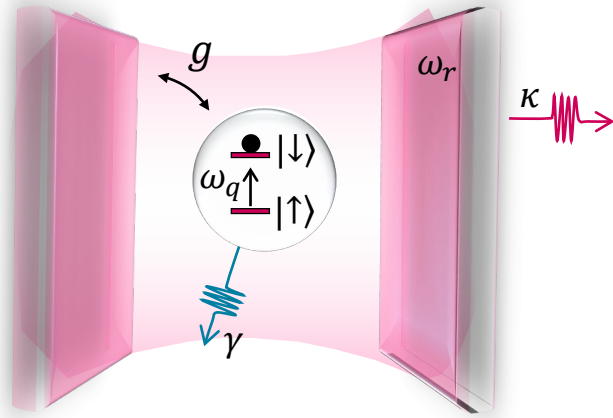
❖ Nanomagnets + **magnonic resonators**

- 😊 GdW10 as spin qubit
- 😊 SCB as cavity
- 😊 Spin – magnon

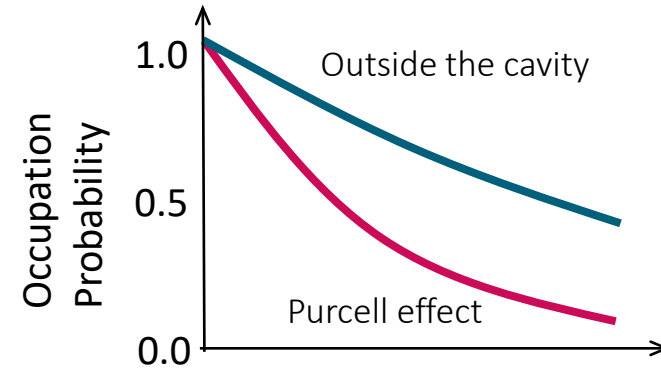
Cavity QED: electromagnetic cavities and qubits



Cavity QED: electromagnetic cavities and qubits



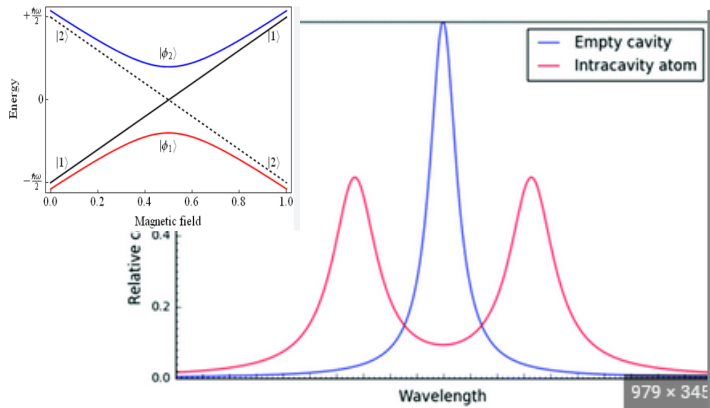
Weak coupling: $\omega_r = \omega_q$



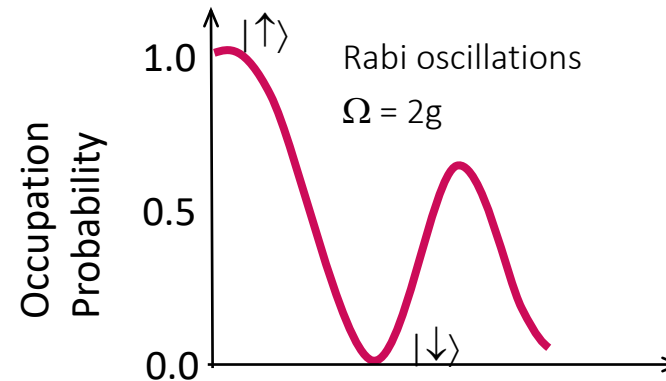
$$g < \gamma, \kappa$$

The atom decay rate increases:

$$\tilde{\gamma} = \gamma + \frac{g^2}{\kappa}$$



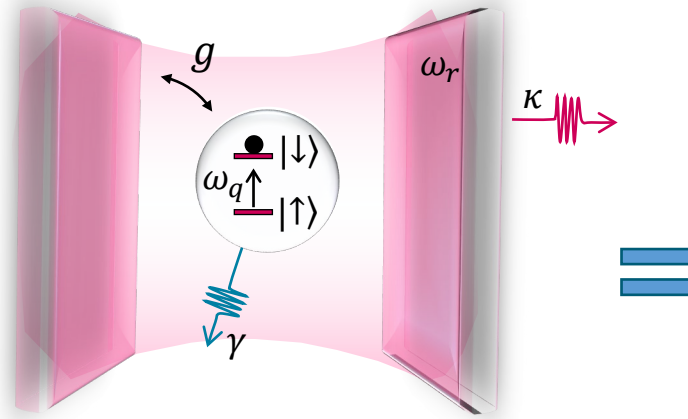
Strong coupling: $\omega_r = \omega_q$



$$g > \gamma, \kappa$$

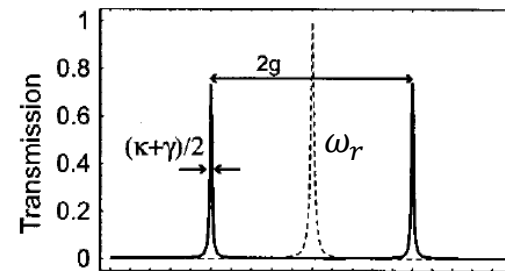
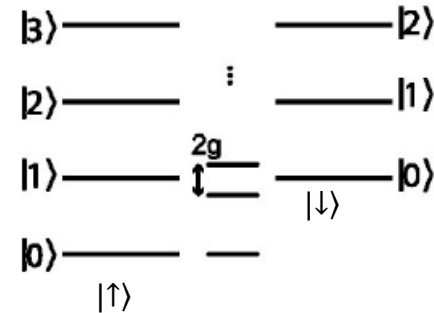
Avoided crossing at degeneracy point

$$\omega_{\pm} = \frac{\omega_r + \omega_q}{2} \pm \frac{\sqrt{4g^2 + (\omega_r - \omega_q)^2}}{2}$$



In resonance:

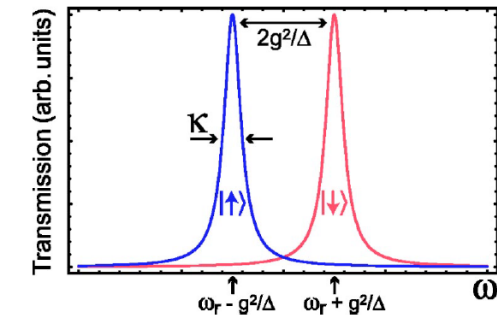
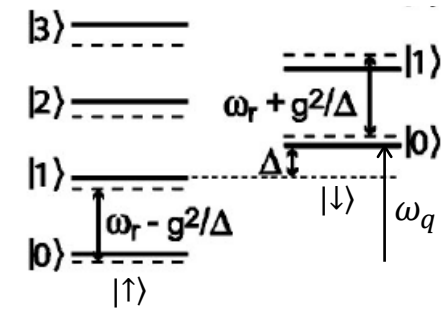
$$\Delta = 0$$



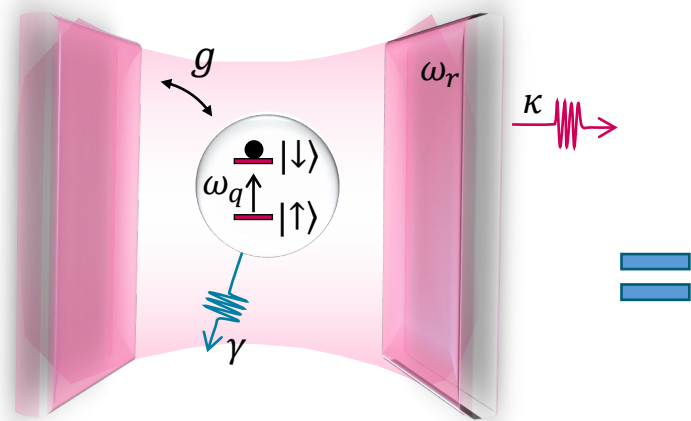
Blais et al PRA 69 2004

Far from resonance:

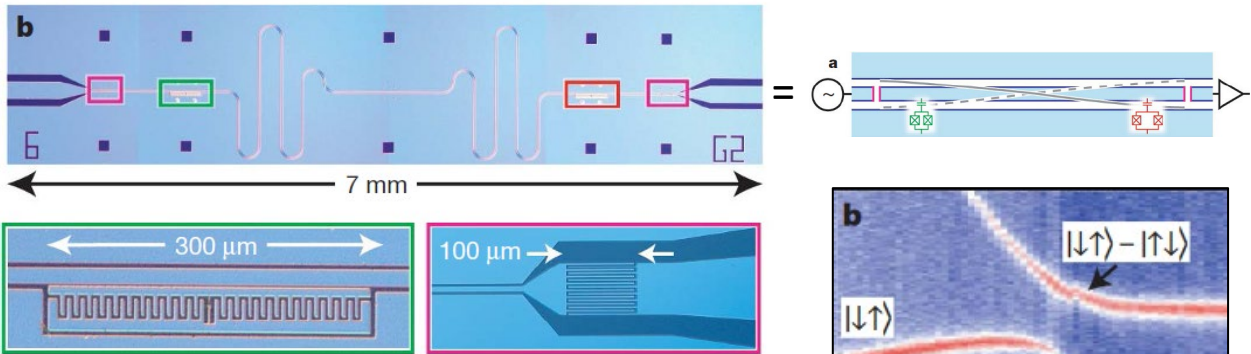
$$\Delta = \omega_q - \omega_r$$



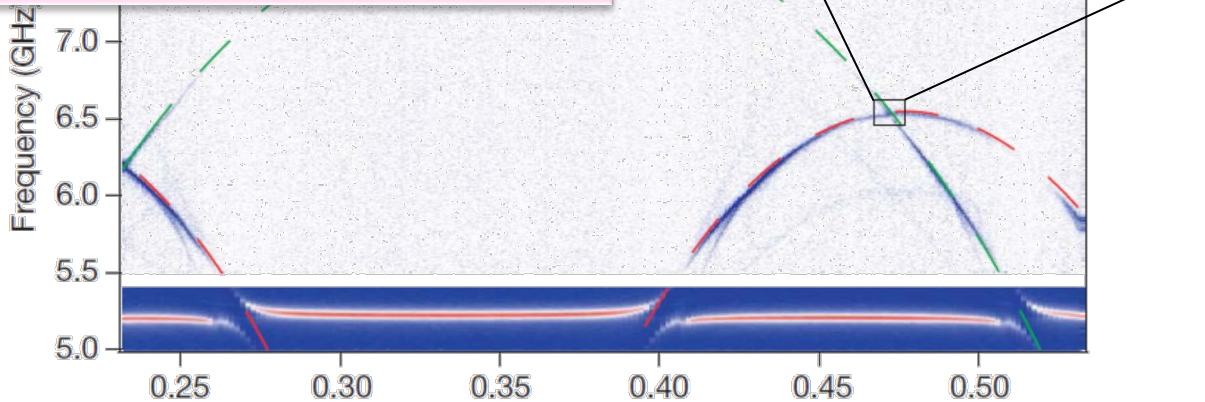
Dispersive READOUT of QUBIT states and PHOTON COUNT!



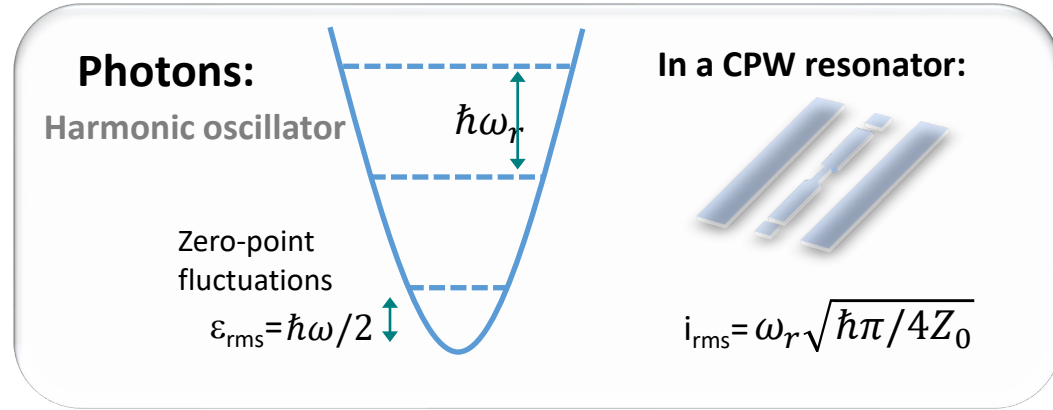
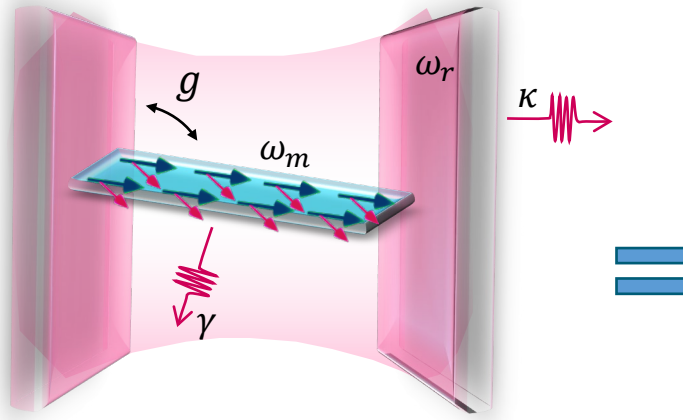
Coupling superconducting qubits via a cavity bus
J. Majer^{1*}, J. M. Chow^{1*}, J. M. Gambetta¹, Jens Koch¹, B. R. Johnson¹, J. A. Schreier¹, L. Frunzio¹, D. I. Schuster¹, A. A. Houck¹, A. Wallraff², A. Blais¹, M. H. Devoret¹, S. M. Girvin¹ & R. J. Schoelkopf²



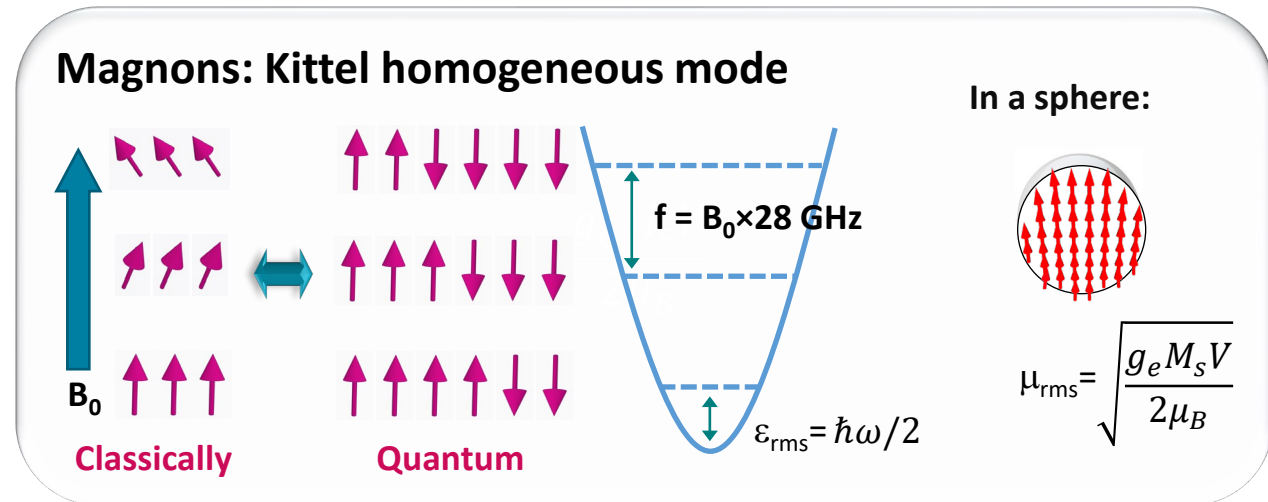
Mediate interaction between distant QUBITS



Cavity QED with bosonic excitations: two coupled harmonic oscillators

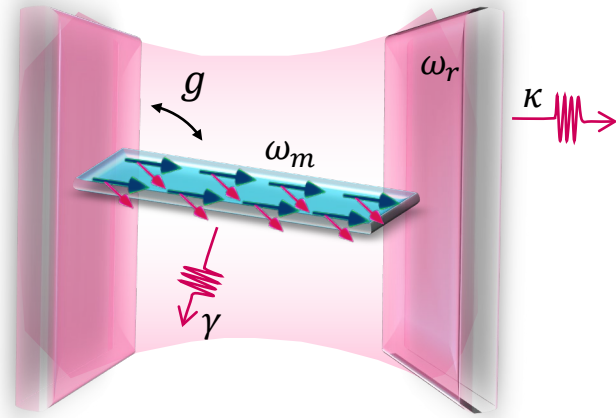


+

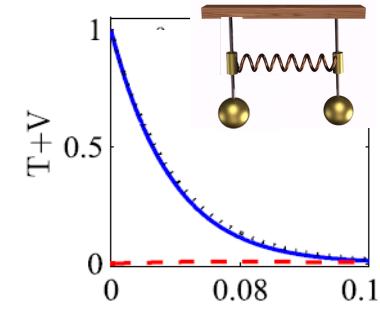
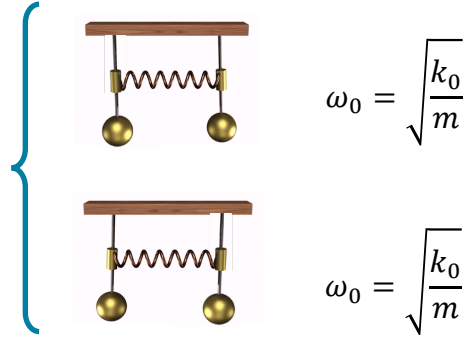
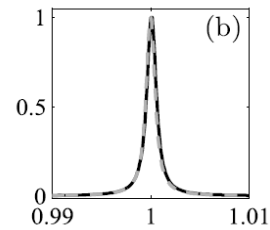


Cavity QED with bosonic excitations: two coupled harmonic oscillators

Two coupled harmonic oscillators with detuning Δk



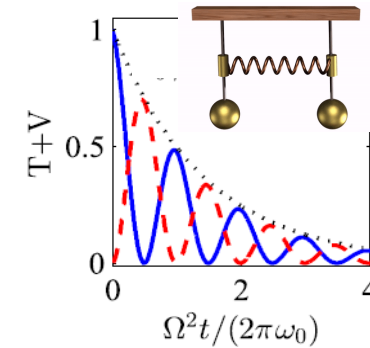
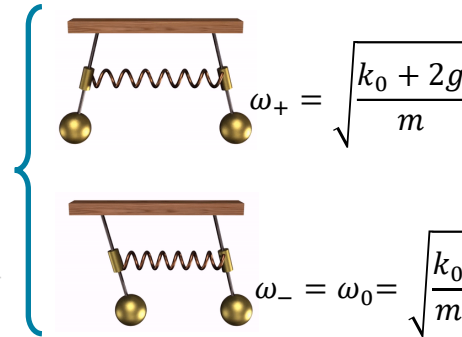
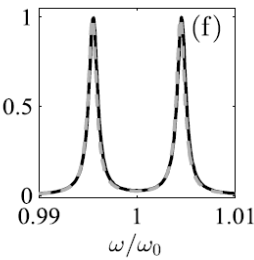
Weak coupling with no detuning $\Delta k=0$



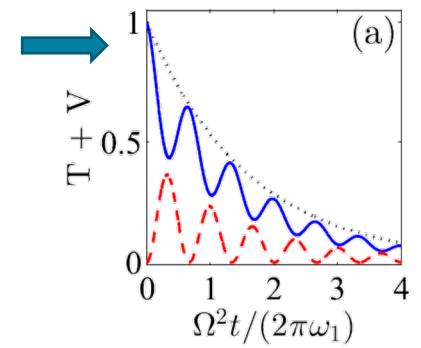
Equivalent to purcell effect!!

$$\kappa = \tilde{\kappa} + \frac{\gamma g^2}{(\omega_p - \omega_m)^2 + \gamma^2}$$

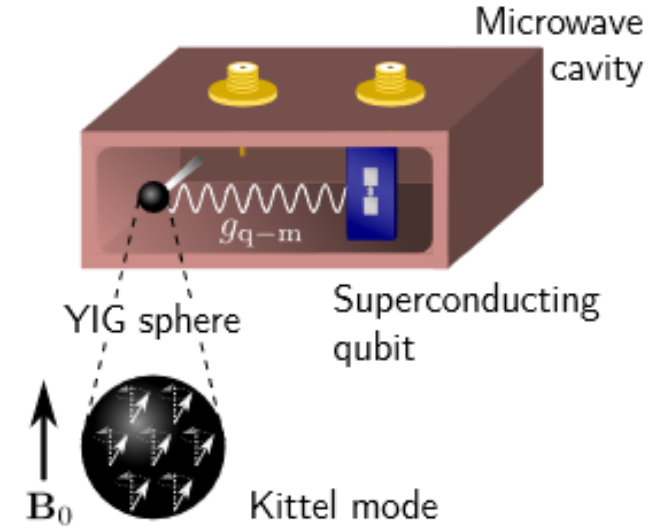
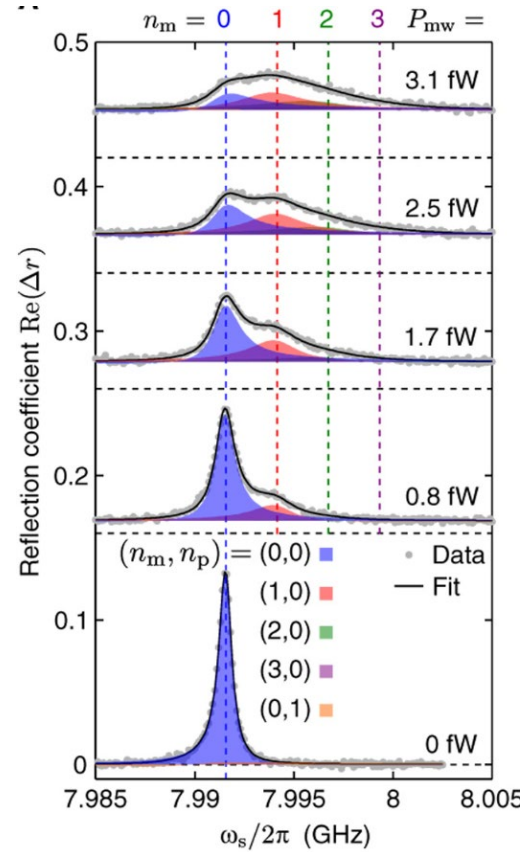
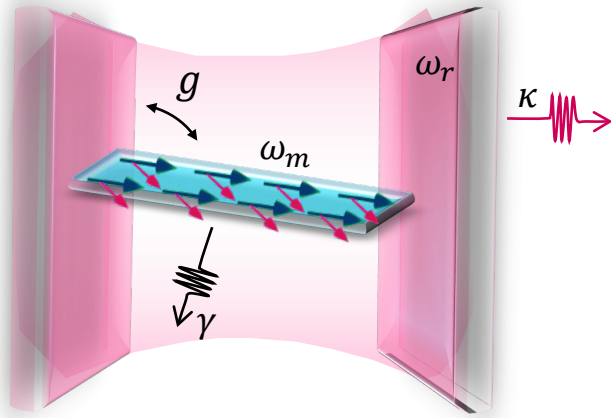
Strong coupling with no detuning $\Delta k=0$



Even detuned $\Delta k \neq 0$



Cavity QED with bosonic excitations: two coupled harmonic oscillators



D. Lachance-Quirion et al, Science Advances 3 (2017)

Counting individual
MAGNONS and **SENSING!**

❖ Light-matter interaction and circuit QED

❖ Nanomagnets + **superconducting resonators**

😊 Coupling to spins

😊 Coupling to magnons

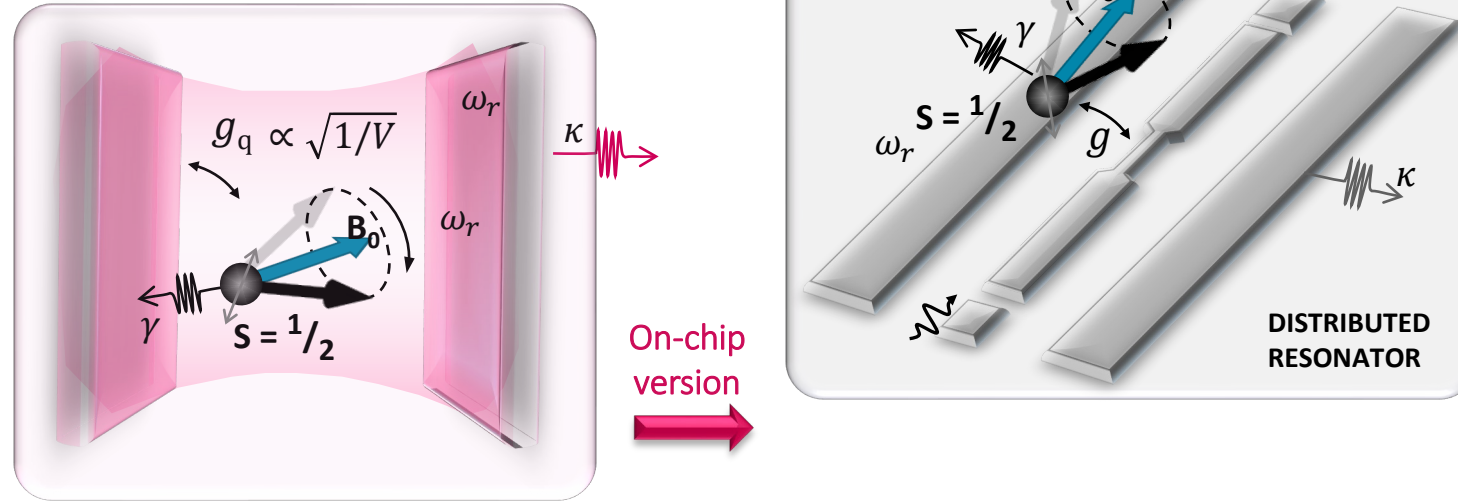
❖ Nanomagnets + **magnonic resonators**

😊 GdW10 as spin qubit

😊 SCB as cavity

😊 Spin – magnon

How can we increase the coupling: From superconducting → to magnonic circuits



$$g = |\langle \uparrow | \boldsymbol{\mu} \mathbf{b}_{\text{rms}} | \downarrow \rangle| \propto b_{\text{rms}} = \sqrt{\frac{\mu_0 \hbar \omega_r}{2V}}$$

$$\frac{1}{4} \hbar \omega = \int_V \frac{1}{2} \frac{b_{\text{rms}}^2}{\mu_0} dV = \frac{1}{2} \frac{L i^2}{2}$$

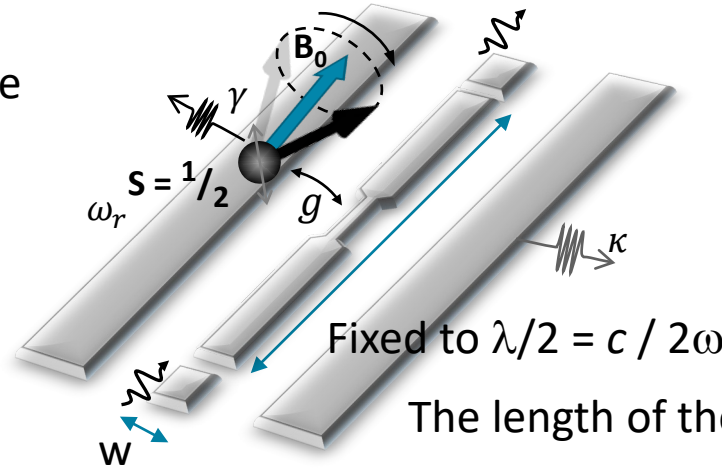
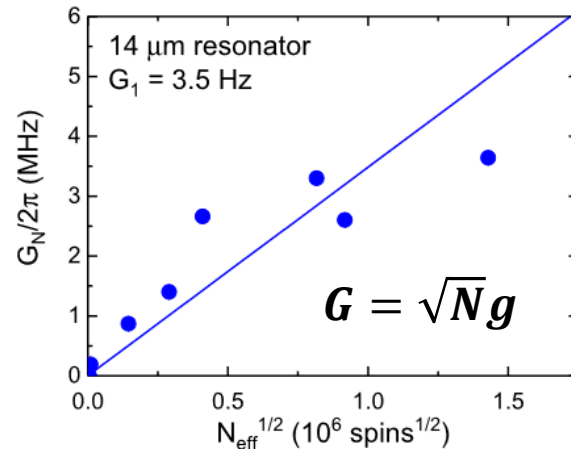
- ❖ Light-matter interaction and circuit QED
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Spin – Photon coupling: the problem to couple to spin qubits

Coupling can be increased by increasing the number of qubits

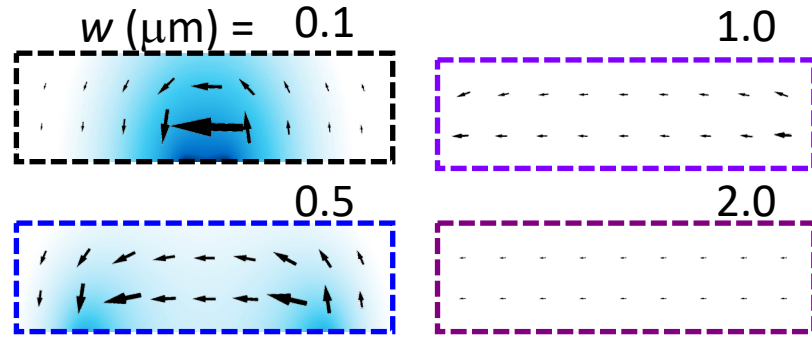
$$G = \sqrt{\sum_i^N g_i^2}$$

Experimental demonstration

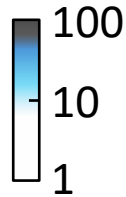


$$g \propto b_{\text{rms}} = \sqrt{\frac{\mu_0 \hbar \omega_r}{2V}}$$

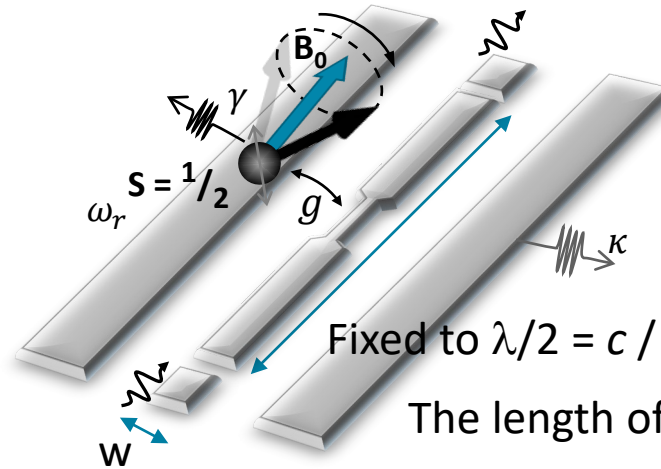
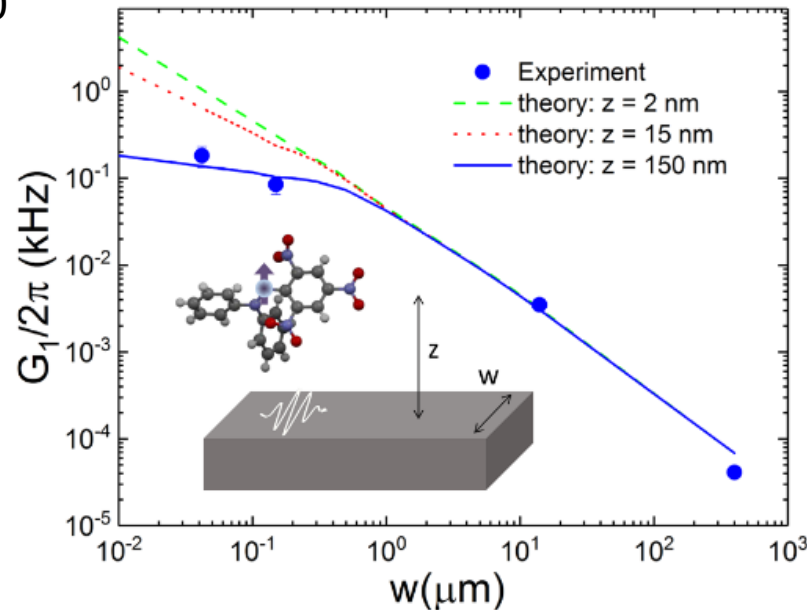
Spin – Photon coupling: the problem to couple to spin qubits



B_{rms} (nT)



Experimental demonstration



The length of the resonator is fixed by frequency

The width can be decreased reducing the mode volume

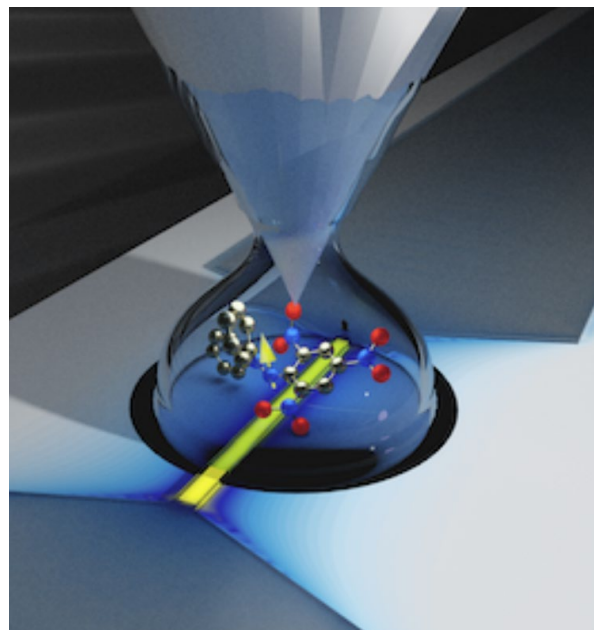
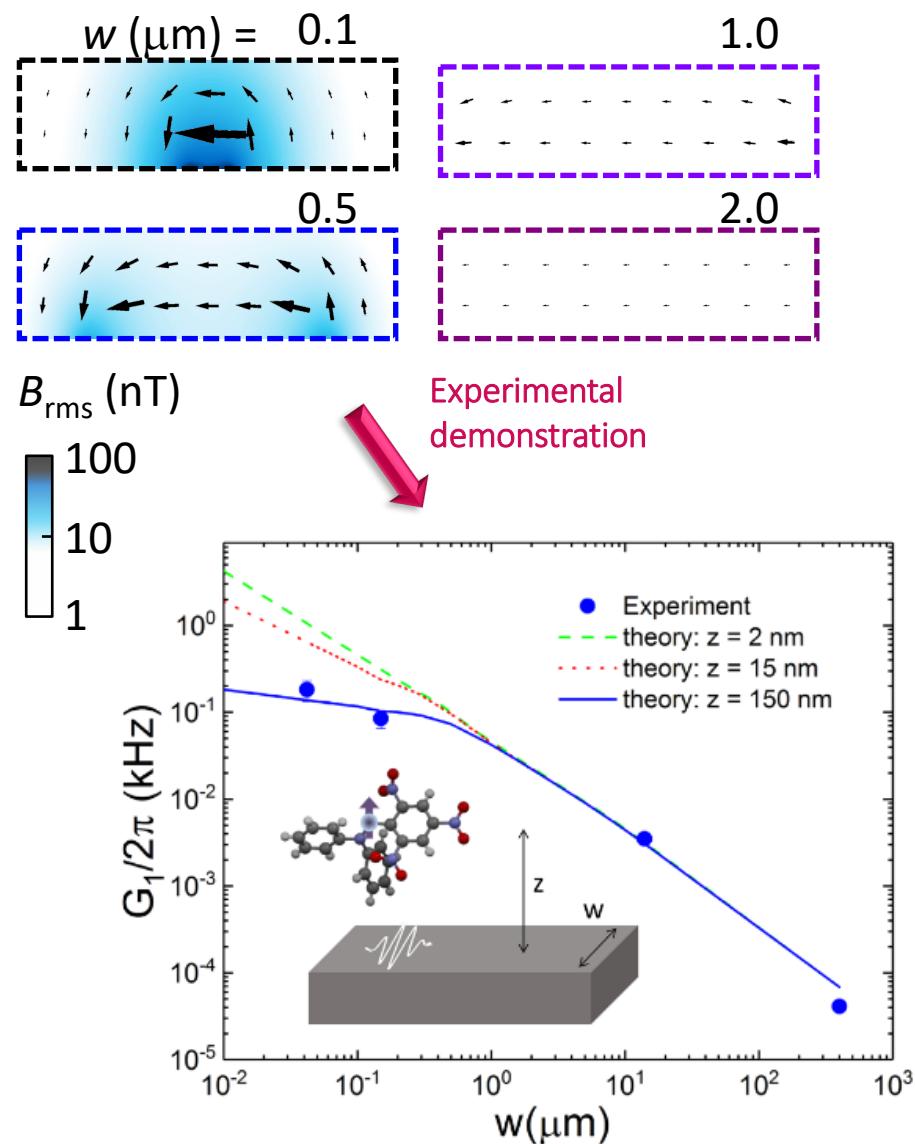
$$g \propto b_{\text{rms}} = \sqrt{\frac{\mu_0 \hbar \omega_r}{2V}} \sim \frac{\omega_r}{r} \sqrt{\frac{2\mu_0}{\pi c}}$$

How can we estimate the z.p.f. field

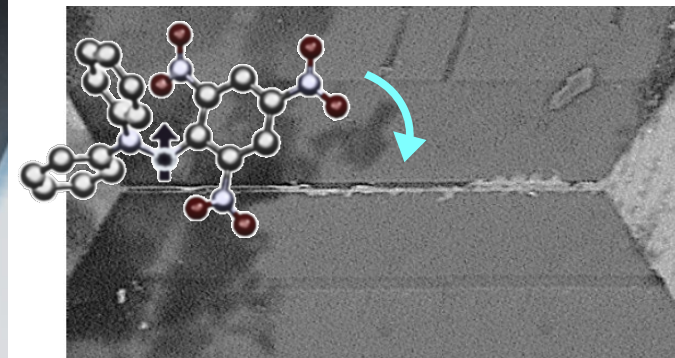
$$\frac{1}{4} \hbar \omega = \int_V \frac{1}{2} \frac{b_{\text{rms}}^2}{\mu_0} dV = \frac{1}{2} \frac{L i_{\text{rms}}^2}{2} \quad \rightarrow \quad i_{\text{rms}} = \omega_r \sqrt{\frac{\hbar \pi}{4Z_0}}$$

$Z_0 = \frac{\pi}{2} L \omega_r$

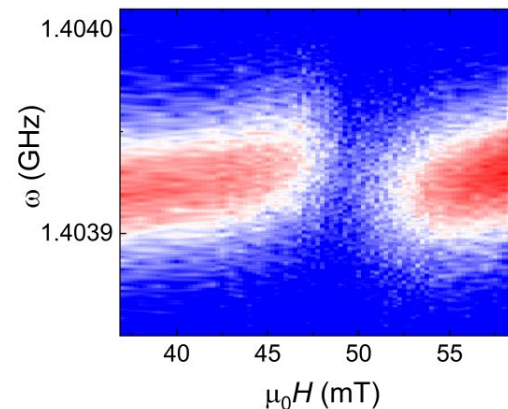
Spin – Photon coupling: the problem to couple to spin qubits



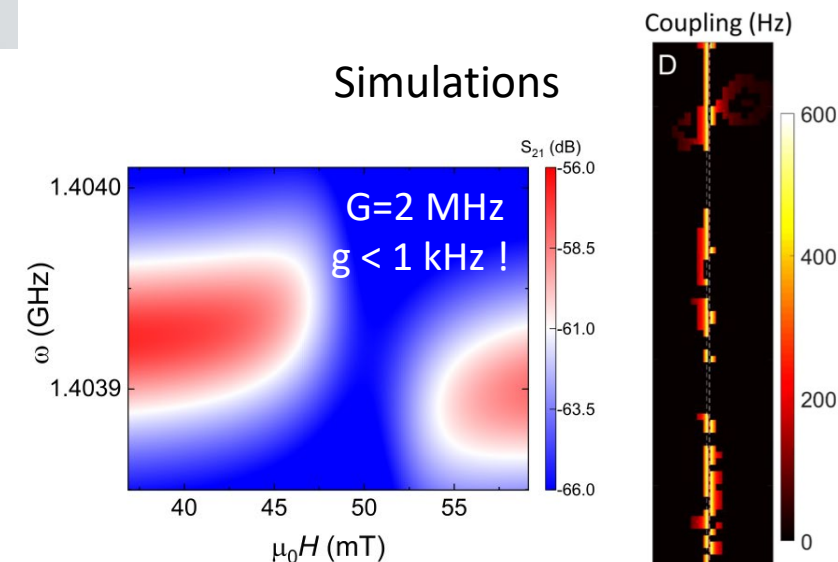
Dip pen distribution of $s=1/2$ molecules (DPPH)



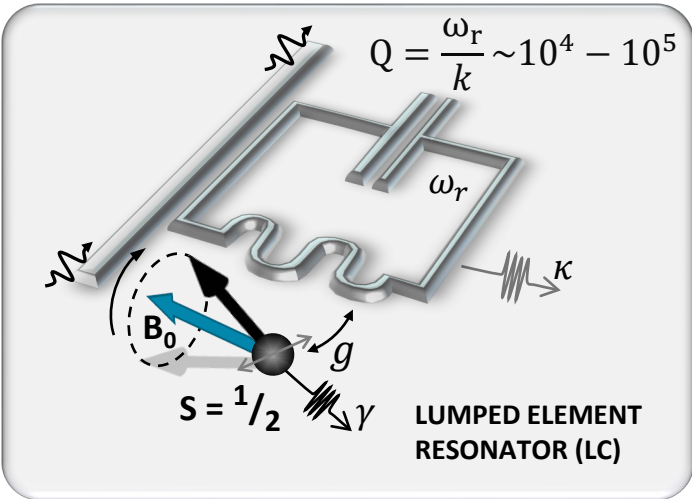
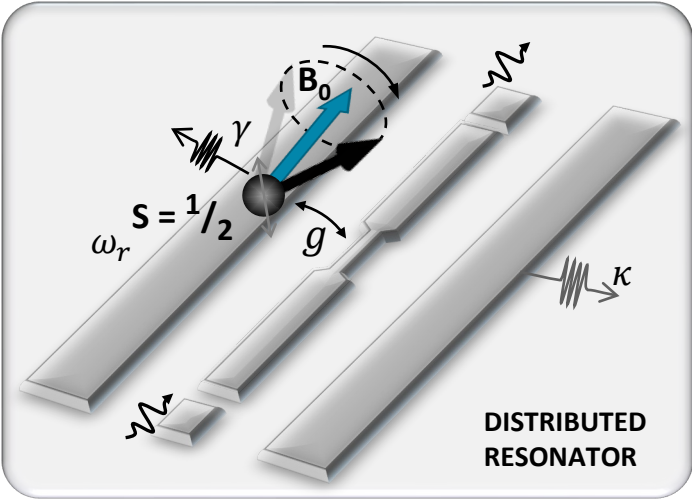
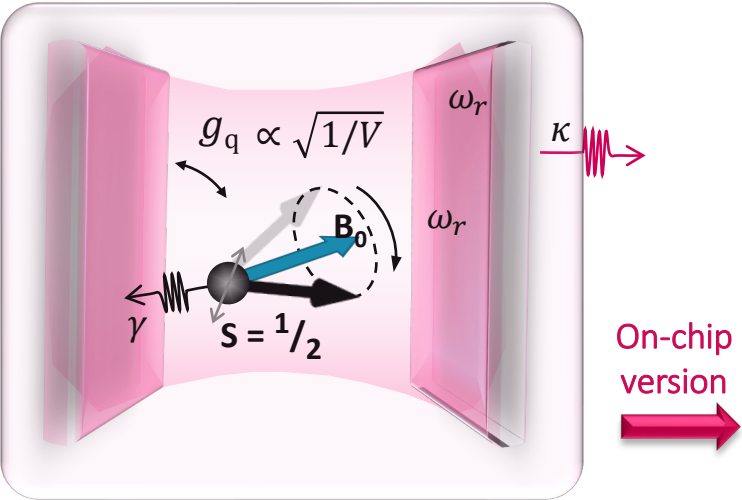
Experiment



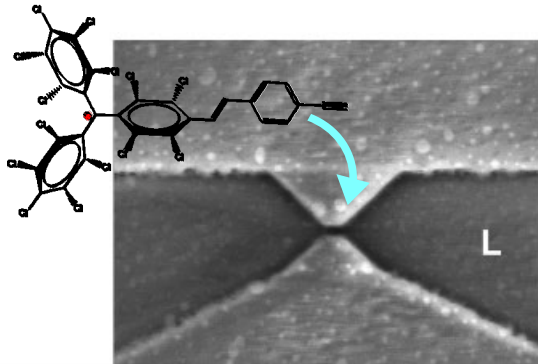
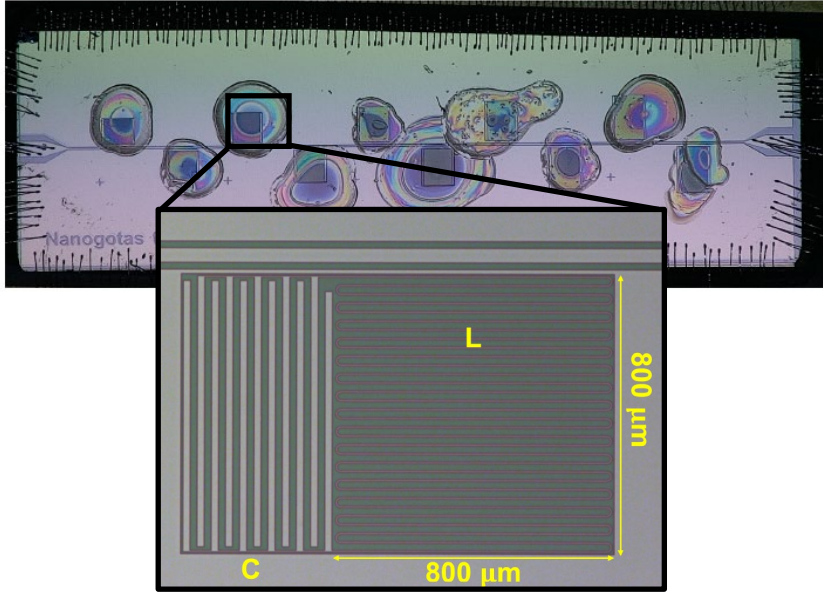
Simulations



How can we increase the coupling: From superconducting → to magnonic circuits



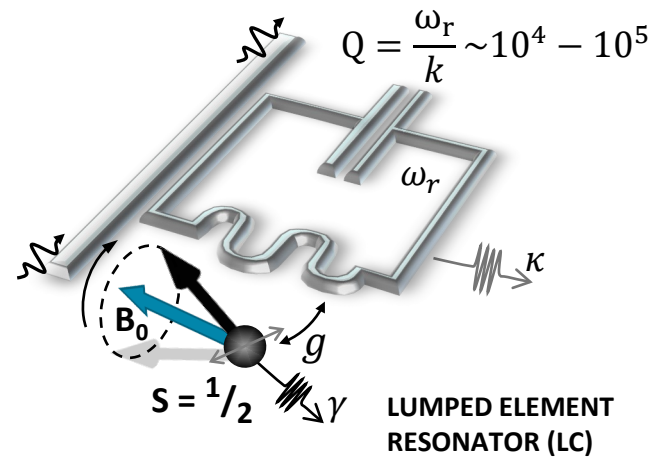
How can we increase the coupling: From superconducting → to magnonic circuits



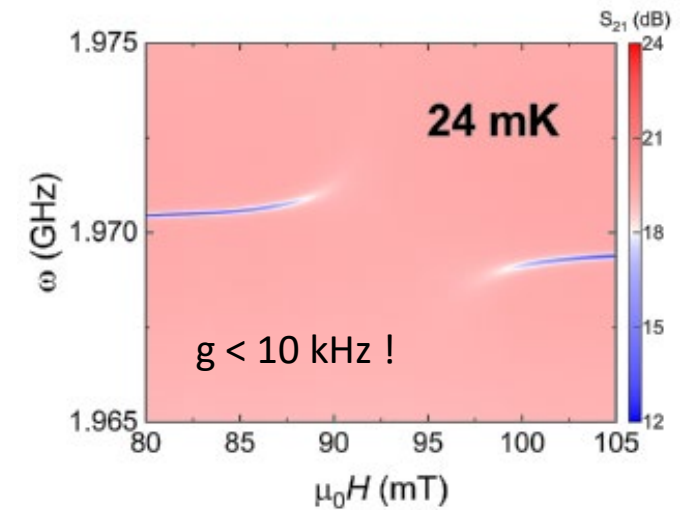
We are still far from coupling to one individual spin qubit (we need MHz)

How can we estimate the z.p.f. field

$$\frac{1}{4} \hbar \omega = \int_V \frac{1}{2} \frac{b_{\text{rms}}^2}{\mu_0} dV = \frac{1}{2} \frac{L i_{\text{rms}}^2}{2} \quad \xrightarrow{Z_0 = L \omega_r} \quad i_{\text{rms}} = \omega_r \sqrt{\frac{\hbar}{Z_0}} = \sqrt{\frac{\hbar \omega_r}{L}}$$



Coupling increases by one order of magnitude !!!



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Magnon nanocavities: how can we estimate the magnon-photon coupling

The hamiltonian, harmonic magnon mode:

$$H = \underbrace{\hbar\omega_c a_c^\dagger a_c}_{H_{\text{cavity}}} + \underbrace{\hbar\omega_v a_m^\dagger a_m}_{H_{\text{vortex}}} + \underbrace{\hbar g (a_c^\dagger + a_m)(a_m^\dagger + a_c)}_{H_{\text{coupling}}}$$

Zeeman magnon-cavity coupling: $H_{\text{coupling}} = \sum_j \mu_j B(r_j) = VM_j B(r_j)$

$$\left\{ \begin{array}{l} \text{cavity quant: } \hat{B} = b(a_c^\dagger + a_c) \\ \text{magnon quant: } \hat{M} = m(a_m^\dagger + a_m) \end{array} \right. \rightarrow \hbar g = Vmb \left\{ \begin{array}{l} \text{Classical driving field: } H_{\text{coupling}} = \hbar g 2\alpha \cos \omega t (a_m^\dagger + a_m) \\ \text{Response of the driven HO: } M(t) = m \langle a_m^\dagger + a_m \rangle = \underbrace{\frac{\hbar g}{Vb} \frac{\hbar g 2\alpha}{\hbar \Delta \omega / 2}}_{\Delta M} \sin \omega t \end{array} \right.$$

one photon

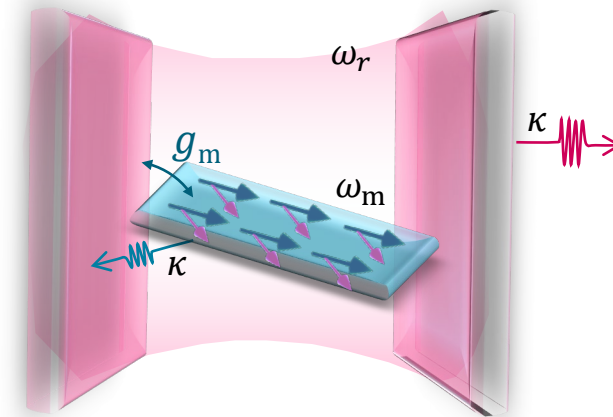


$$g = \frac{b_{\text{rms}}}{2} \sqrt{\frac{V \chi \Delta f_m}{\hbar}}$$

$$\chi(f_m) = \Delta M / b$$

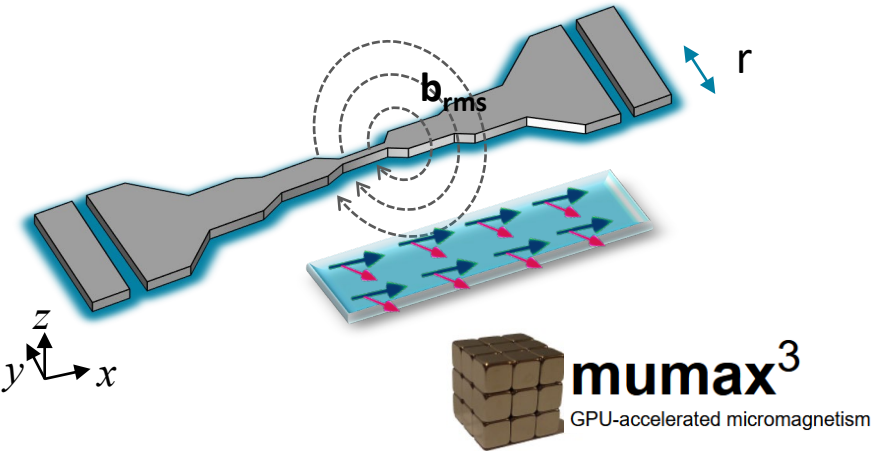
One needs to know the susceptibility of the mode

MJ M-P & D Zueco, ACS photonics 2019



Spin – Photon coupling: Coupling to ferromagnetic spins

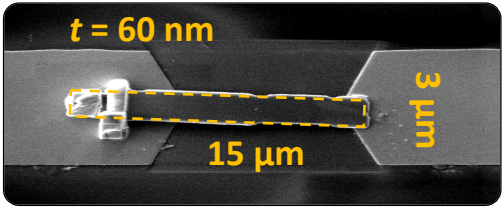
Fixed to $\lambda/2 = c / 2\omega$



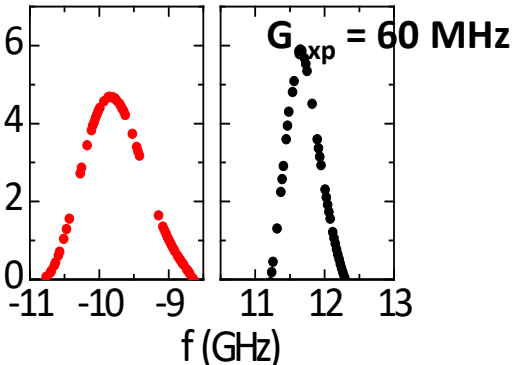
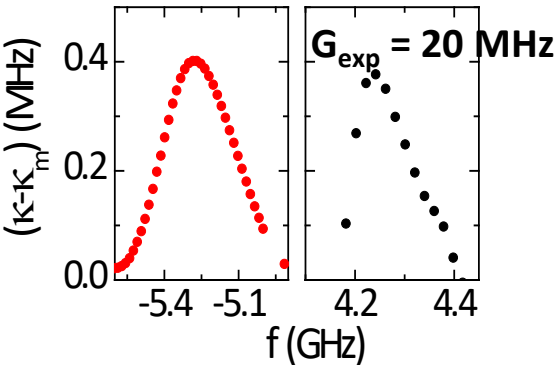
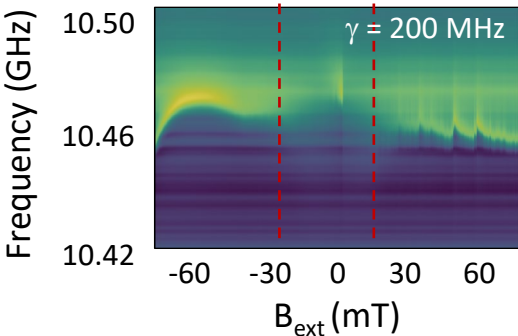
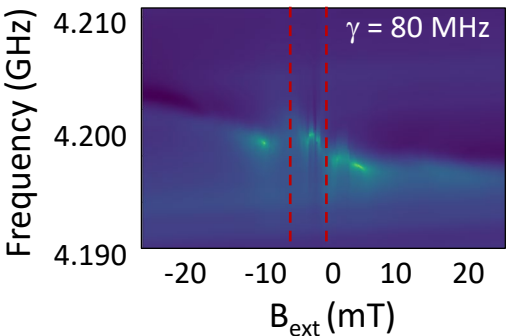
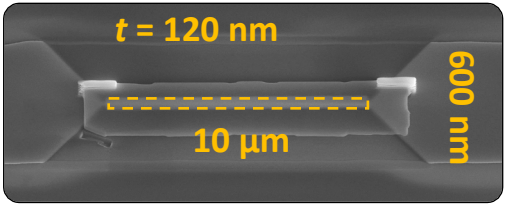
Coupled spins →

$$G = \sqrt{\frac{\gamma v_{\text{cell}}}{4\hbar} \sum_j^{N_{\text{cell}}} \Delta \mathbf{M}_j \cdot \mathbf{b}_{\text{rms}}(r_j)}$$

G_{theo} = 40 MHz



G_{theo} = 80 MHz



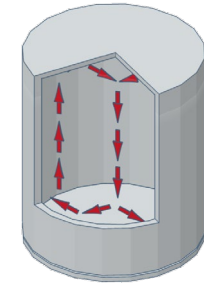
Not in strong coupling

Martinez-Losa, (...) MJ M-P Phys Rev Appl, 2022

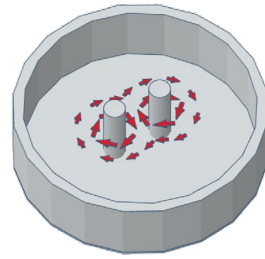
$$\kappa = \tilde{\kappa} + \frac{\gamma g^2}{(\omega_p - \omega_m)^2 + \gamma^2}$$

In the weak coupling regime:

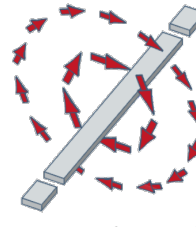
Real cavities yield strongly non-homogeneous profiles...



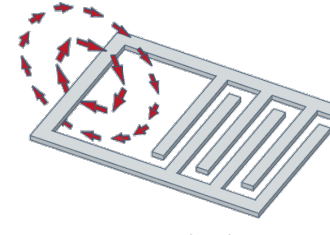
Cylindrical



Re-entrant

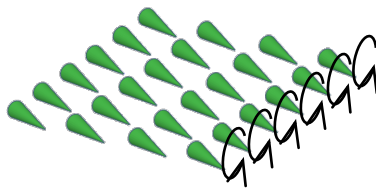


Distributed

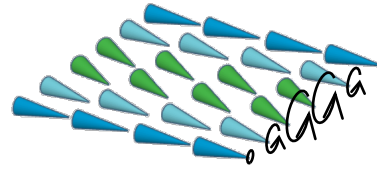


Lumped-element

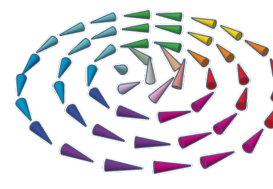
That excite non homogeneous magnon modes. Also magnetic textures can be stabilized



Kittel ($k = 0$)



Magnetostatic ($k \neq 0$)



Vortex



Skyrmion

Heisenberg equation of motion for spin:

$$\dot{\hat{S}}_i = \frac{i}{\hbar} [\hat{H}, \hat{S}_i] \xrightarrow{\text{Large spin limit}} \dot{\mathbf{m}}_i = -\gamma \mathbf{m}_i \times \mathbf{B}_{\text{eff}}(\mathbf{r}_i) + \text{DAMPING}$$

We can now simulate the response of M upon the action of a cavity field B_{rms} . Hamiltonian:

$$\hat{H} = \gamma \sum_i \hat{S}_i \cdot \mathbf{B}_{\text{eff}}(\mathbf{r}_i) + \gamma \sum_i \hat{S}_i \cdot \mathbf{B}_{\text{rms}}(\mathbf{r}_i) (\hat{a} + \hat{a}^\dagger) + \hbar \omega_c \hat{a}^\dagger \hat{a}$$

Corresponding Heisenberg eq. of motion $\rightarrow$
$$\begin{cases} \dot{\hat{S}}_i = -\gamma \hat{S}_i \times \mathbf{B}_{\text{eff}}(\mathbf{r}_i) - \gamma \left(\hat{S}_i \times \mathbf{B}_{\text{rms}}(\mathbf{r}_i) \right) (\hat{a} + \hat{a}^\dagger) , \\ \dot{\hat{a}} = -i\omega_c \hat{a} - i\frac{\gamma}{\hbar} \sum_i \hat{S}_i \cdot \mathbf{B}_{\text{rms}}(\mathbf{r}_i) . \end{cases}$$
 Expectation value $\rightarrow$
$$\begin{cases} \dot{\mathbf{S}}_i = -\gamma \mathbf{S}_i \times \mathbf{B}'_{\text{eff}}(\mathbf{r}_i) , \\ \dot{\alpha} = -i\omega_c \alpha - i\frac{\gamma}{\hbar} \sum_i \mathbf{S}_i \cdot \mathbf{B}_{\text{rms}}(\mathbf{r}_i) , \end{cases}$$

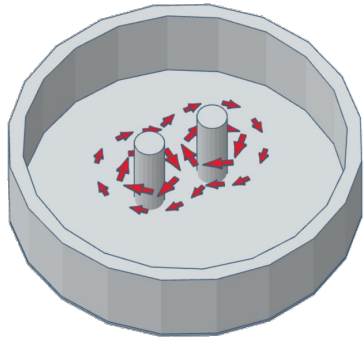
$$\mathbf{B}'_{\text{eff}}(\mathbf{r}_i) = \mathbf{B}_{\text{eff}}(\mathbf{r}_i) + \mathbf{B}_{\text{rms}} \Gamma(t)$$

Memory term

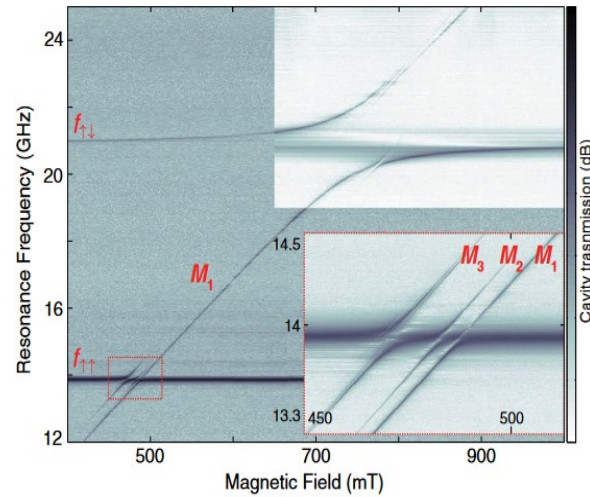
$$\Gamma(t) = -\frac{2M_s V_c}{\hbar} \mathbf{B}_{\text{rms}} \int_0^t e^{-\kappa(\tau-t)} \sin(\omega_c(\tau-t)) \mathbf{m}(\tau) d\tau$$

Martinez-Losa, (...) MJ M-P submitted to
Computer Physics Commun.

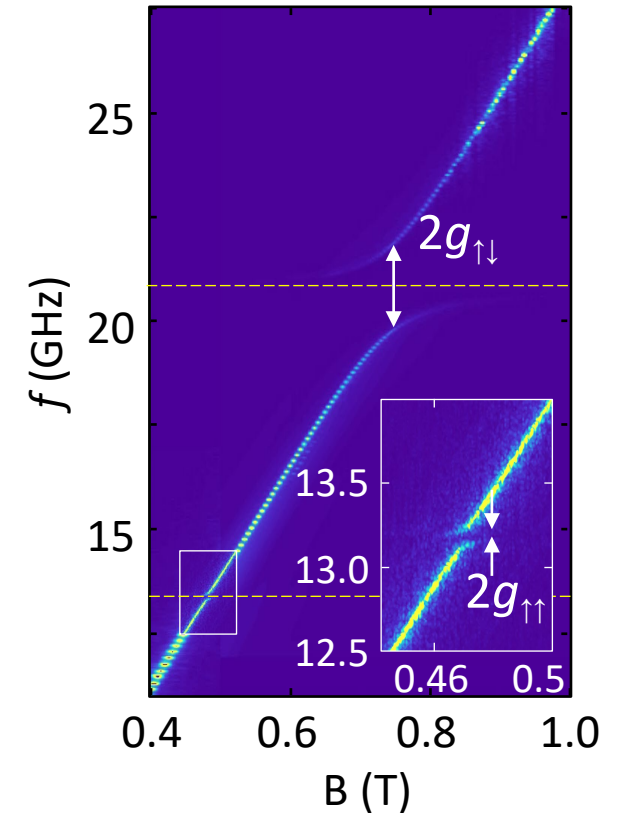
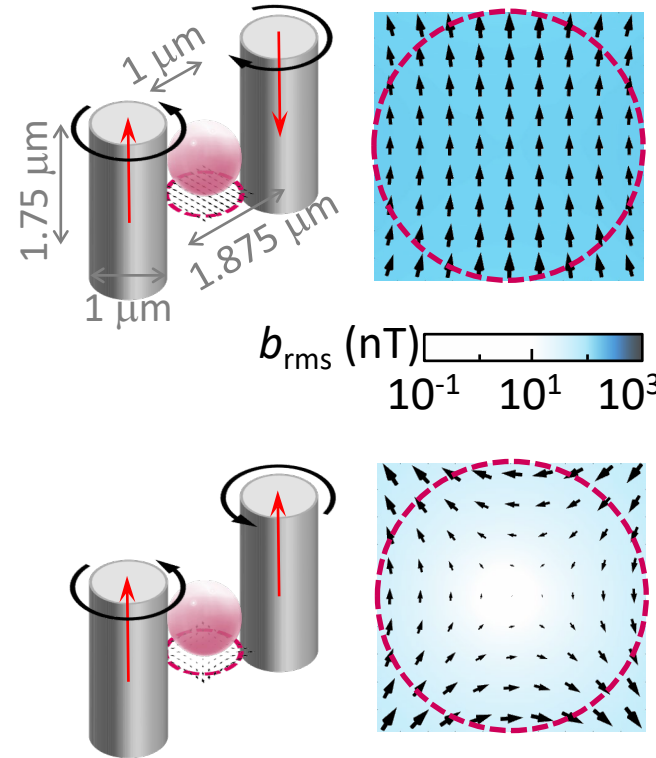
Reproduce EXPERIMENT by:



M. Goryachev,, Physical
Review Applied 2 (5)
(2014)

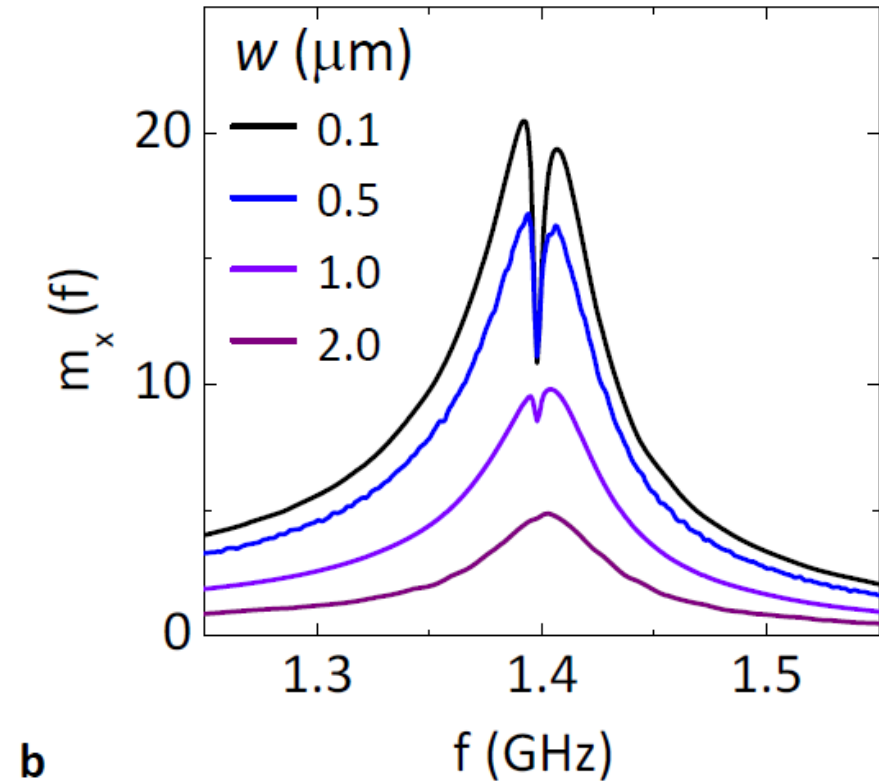
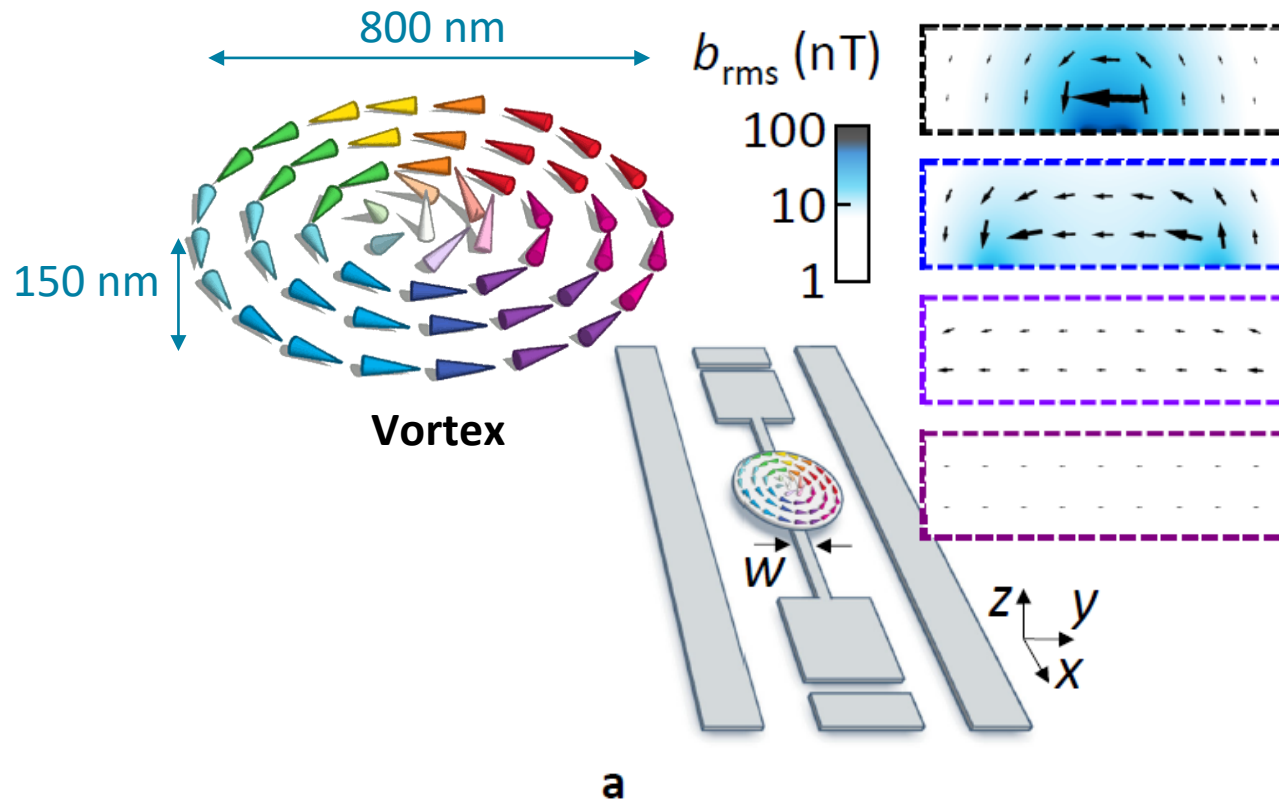


NUMERICAL RESULTS:



Martinez-Losa, (...) MJ M-P SUBMITTED TO
Computer Physics Commun.

Magnetic textures + non-homogeneous magnetic profile



Martinez-Losa, (...) MJ M-P SUBMITTED TO
Computer Physics Commun.

❖ Light-matter interaction and circuit QED

❖ Nanomagnets + **superconducting resonators**

😊 Coupling to spins

😊 Coupling to magnons

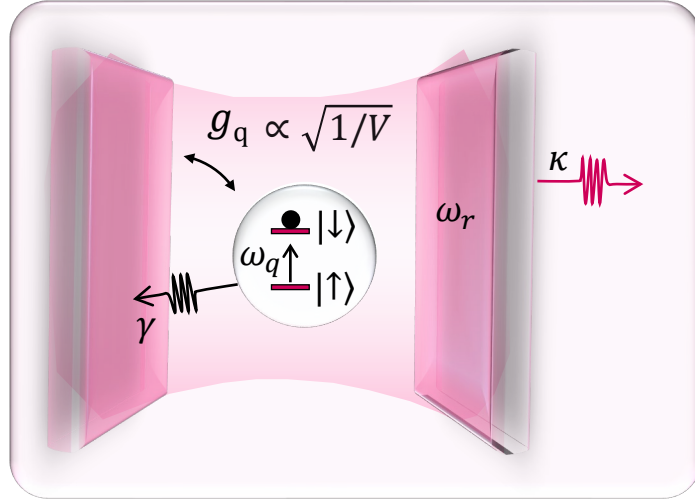
❖ Nanomagnets + **magnonic resonators**

😊 GdW10 as spin qubit

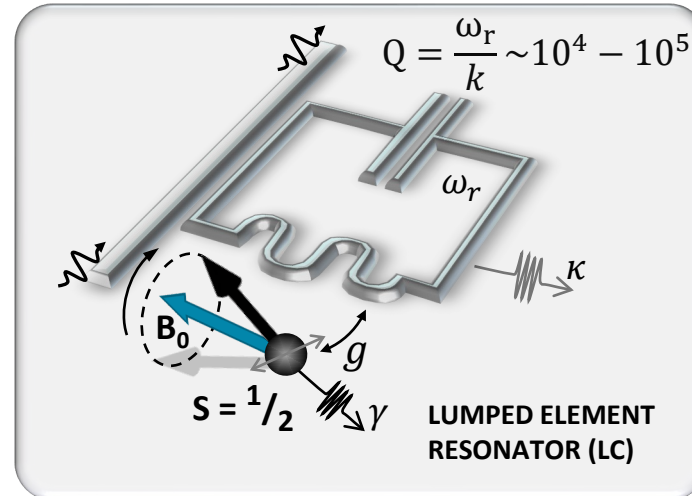
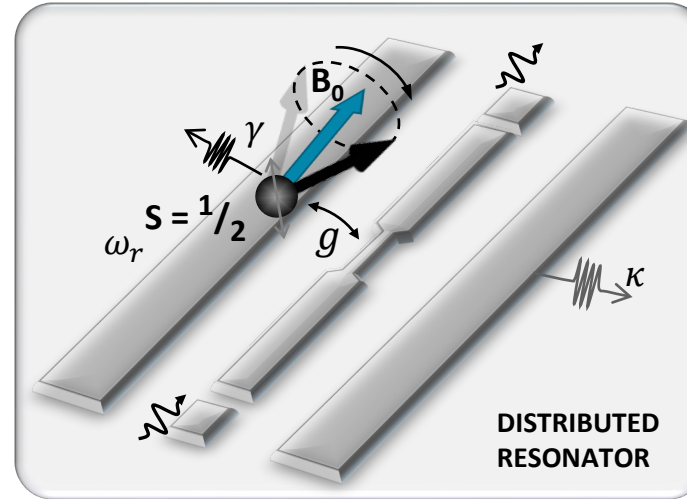
😊 SCB as cavity

😊 Spin – magnon

How can we increase the coupling: From superconducting → to magnonic circuits

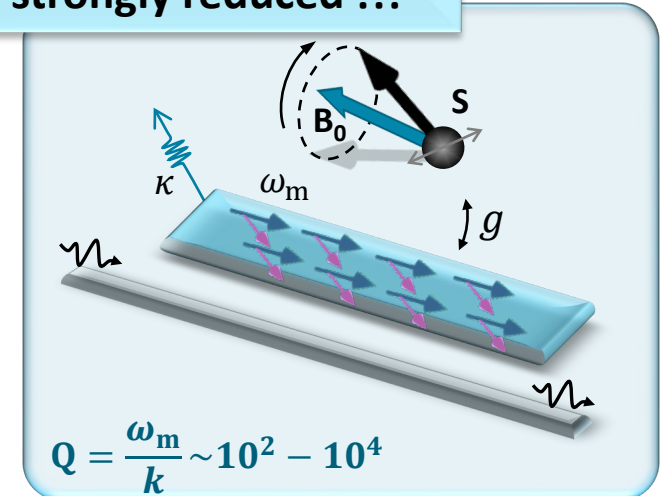


On-chip
version



Magnon
version

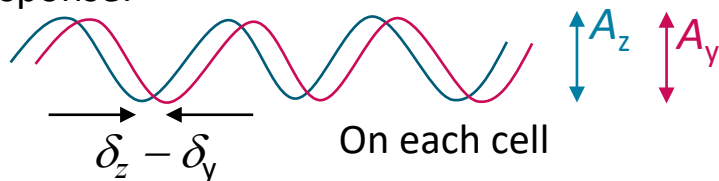
Mode volumen can be
strongly reduced !!!





Use excitation field $B = \beta \sin(\omega t)$

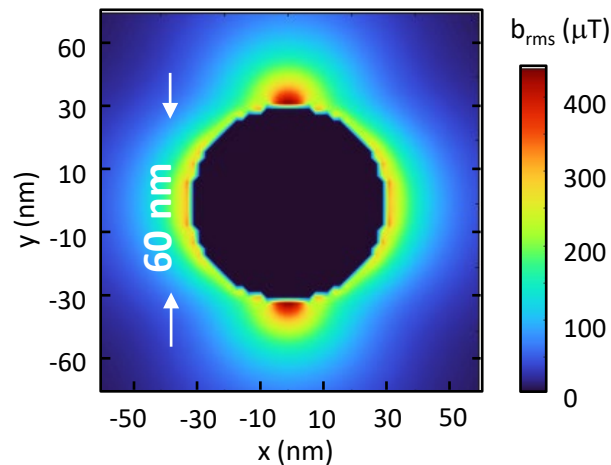
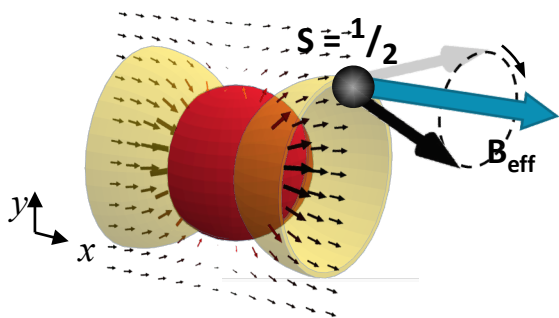
Response:



$$\Lambda(\beta) = \sqrt{\frac{2g_e\mu_B M_z}{v_{\text{cell}} \sum_n A_z(\mathbf{r}_n) A_y(\mathbf{r}_n) |\sin(\delta_z(\mathbf{r}_n) - \delta_y(\mathbf{r}_n))|}}$$

$$\rightarrow \mathbf{B}_{\text{rms}} = \Lambda(\beta) \mathbf{B}_{\text{stray}}^{\text{ac}}(\beta)$$

Ferromagnetic spheres:

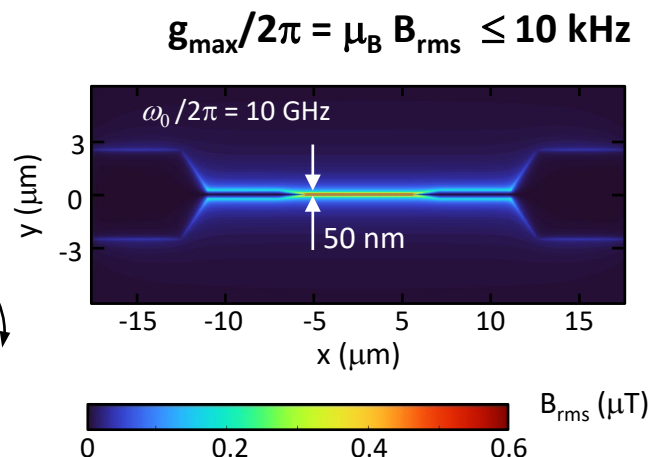
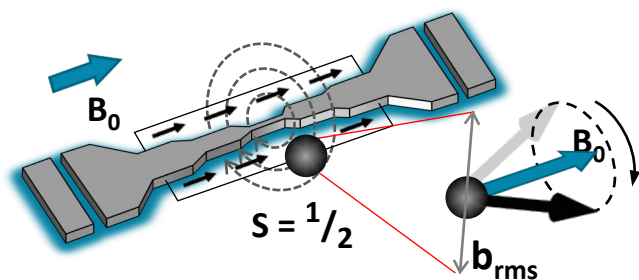


Relevant zpf to calculate coupling to $\frac{1}{2}$ spin:

$$B_{\text{rms}} = \sqrt{B_{\text{rms},1}^2 + B_{\text{rms},2}^2}$$

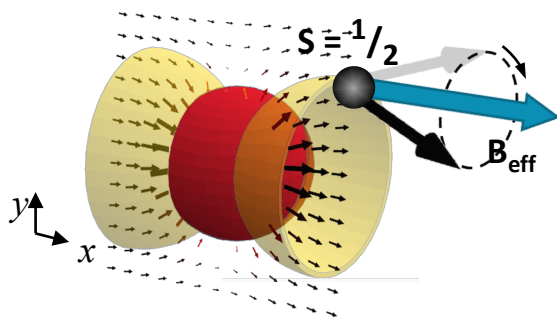
Magnon nanocavities

Superconducting cavities:

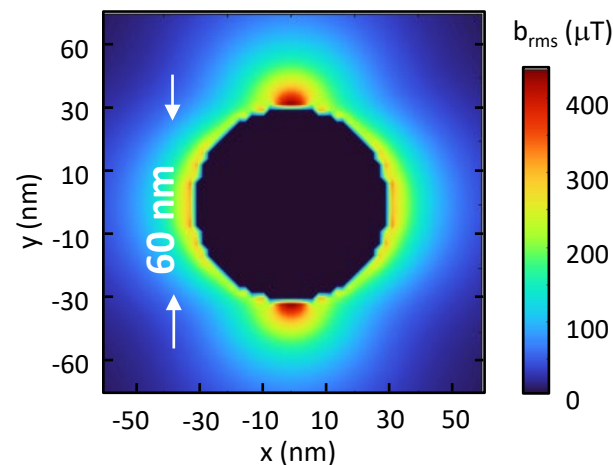


$$g \propto b_{\text{rms}} = \sqrt{\frac{\mu_0 \hbar \omega_r}{2V}}$$

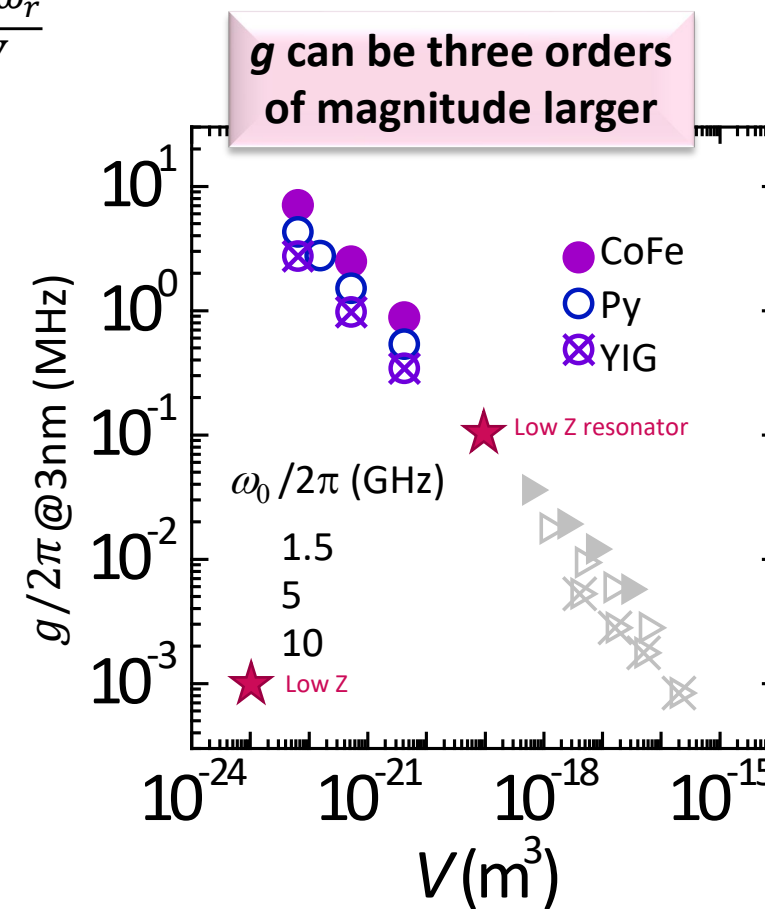
Ferromagnetic spheres:



$g_{\max}/2\pi = \mu_B B_{\text{rms}} \sim \text{MHz for } S=1/2 !!!$

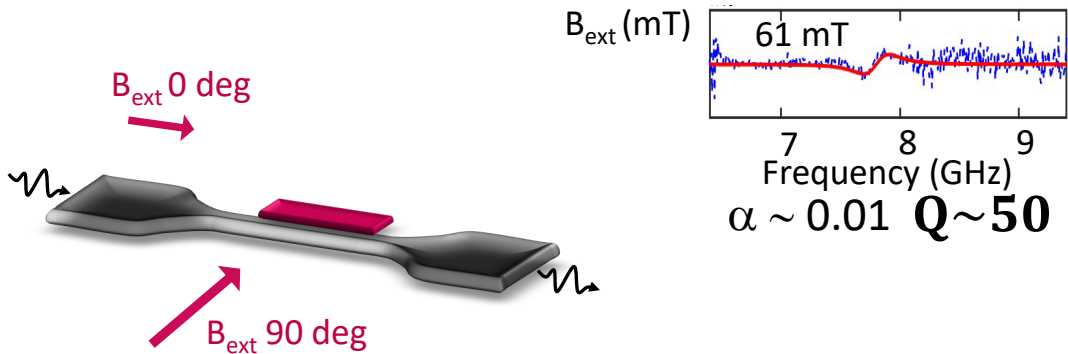
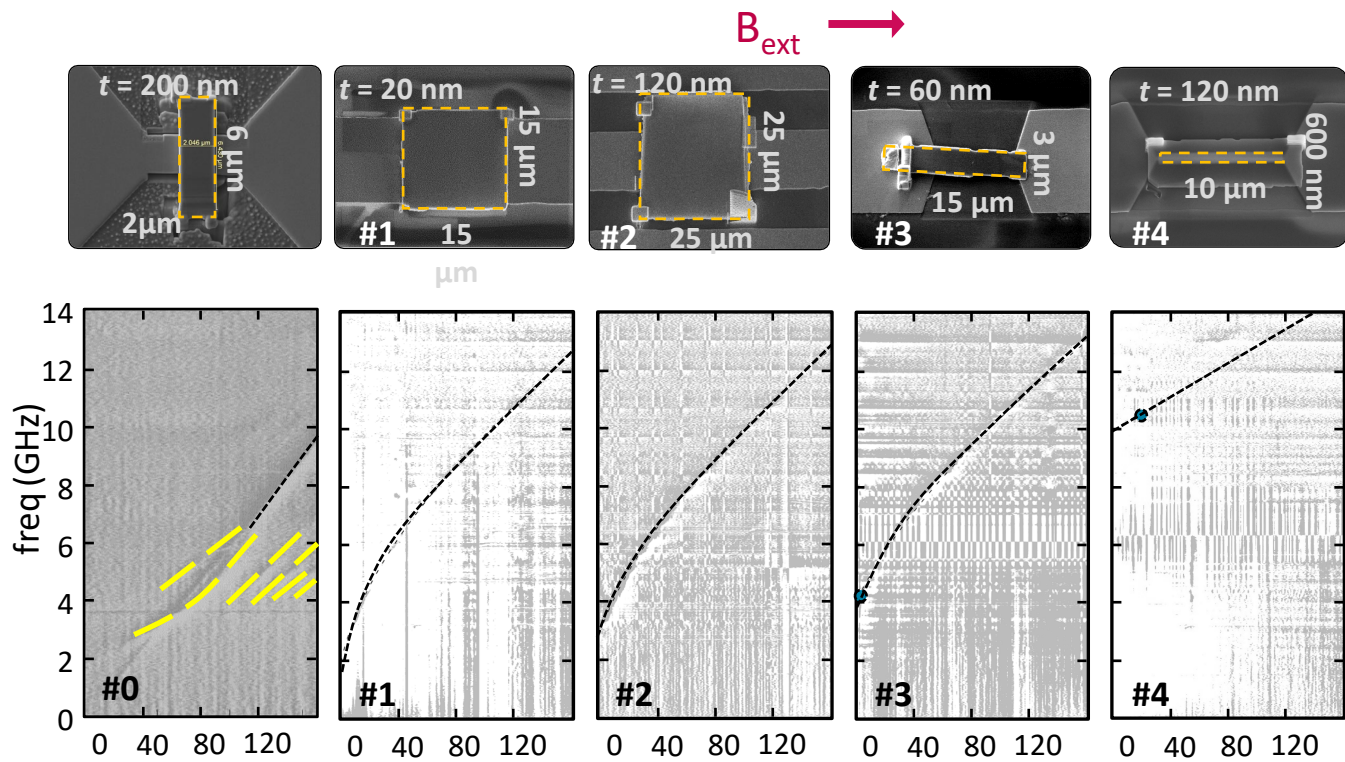


$$g \propto M_{\text{rms}} = \sqrt{\frac{g_e M_s}{2\mu_B V}}$$



Metallic ferromagnets: permalloy and FeCo

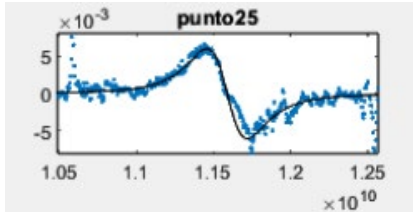
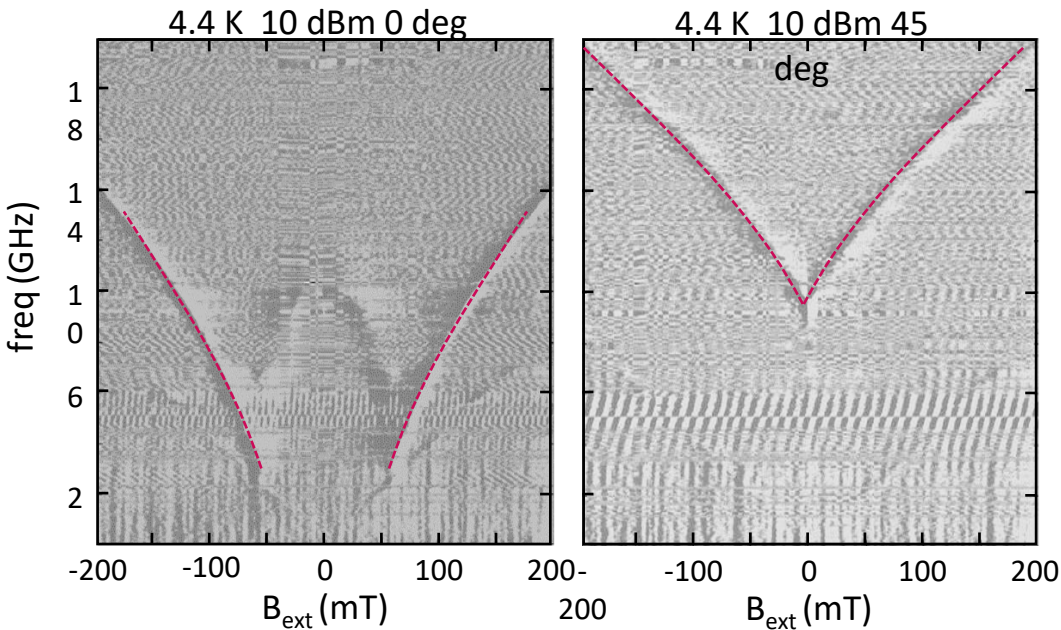
Broad band ferromagnetic resonance at mK temperatures



(collaboration with A. Butera (Bariloche))

FeCo 110 -hard axis

FeCo 100 -easy axis



Shape and magnetocrystalline anisotropy allows to tune the resonance frequency !!

Insulating ferrimagnets: Yttrium Iron Garnet at mK temperatures with no GGG

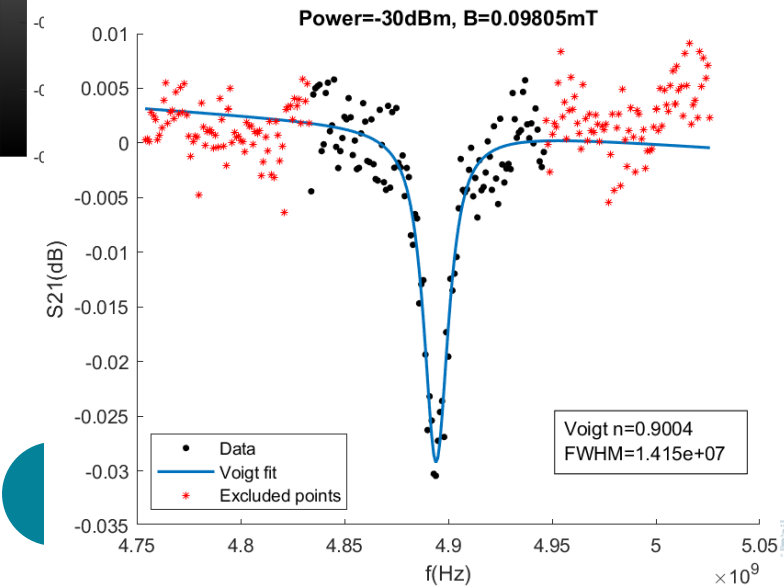
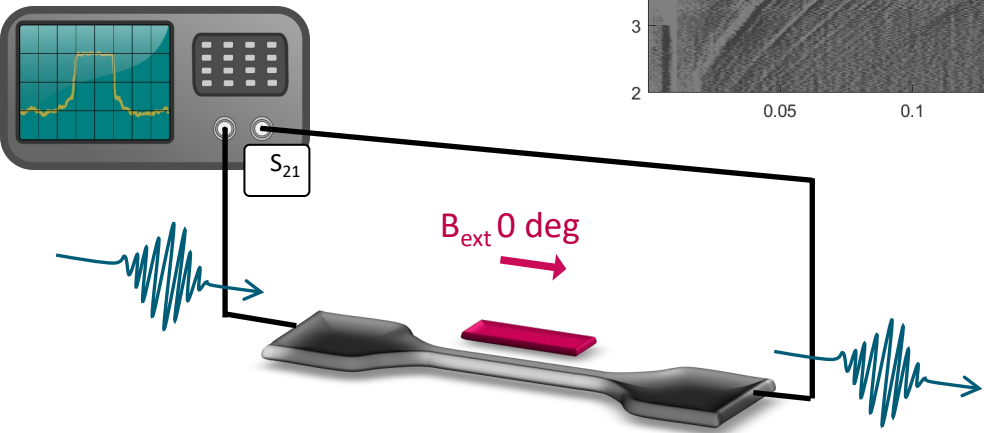
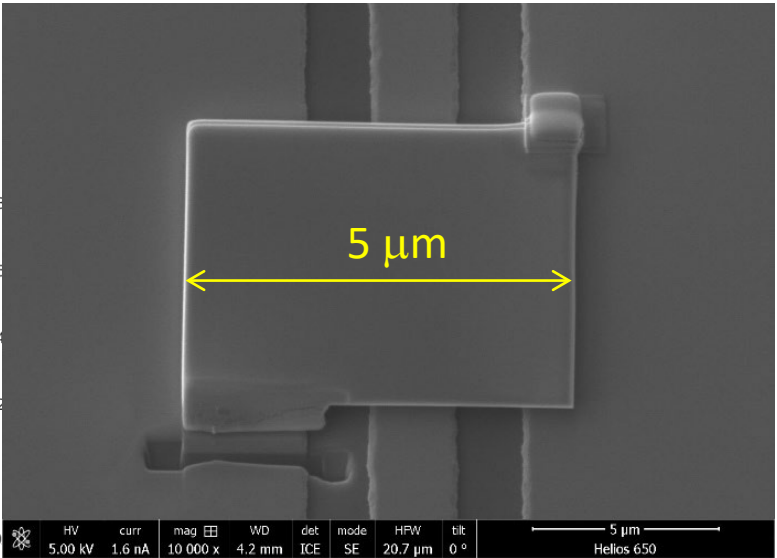
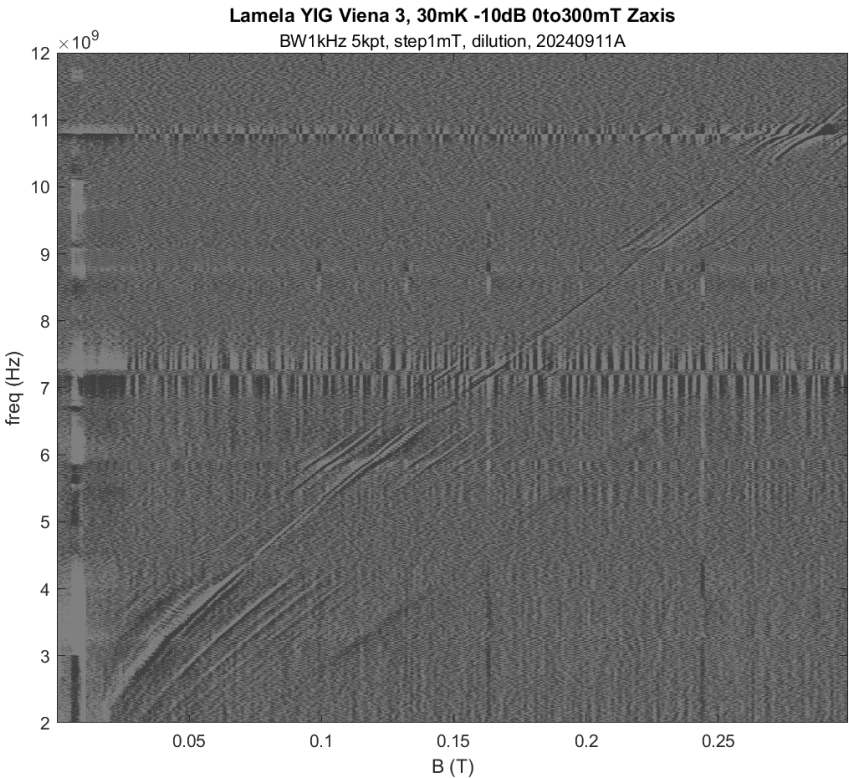
Using YIG at Room temperature: $\alpha \sim 0.0001$ $Q \sim 5000$

Using nanopatterned YIG at 30 mK:

$\alpha \sim 0.0007$

$Q \sim 700$

Very encouraging !!

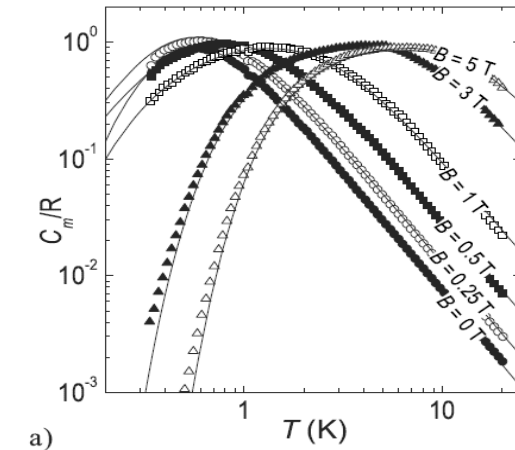
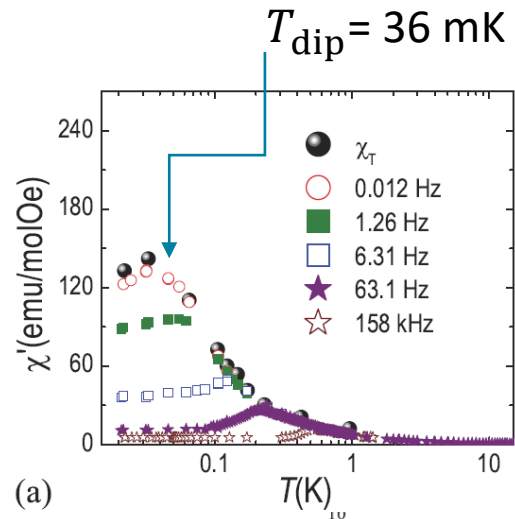
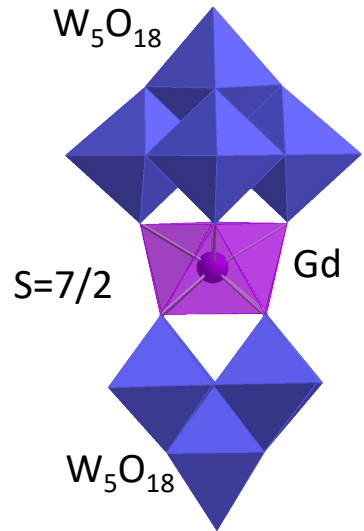


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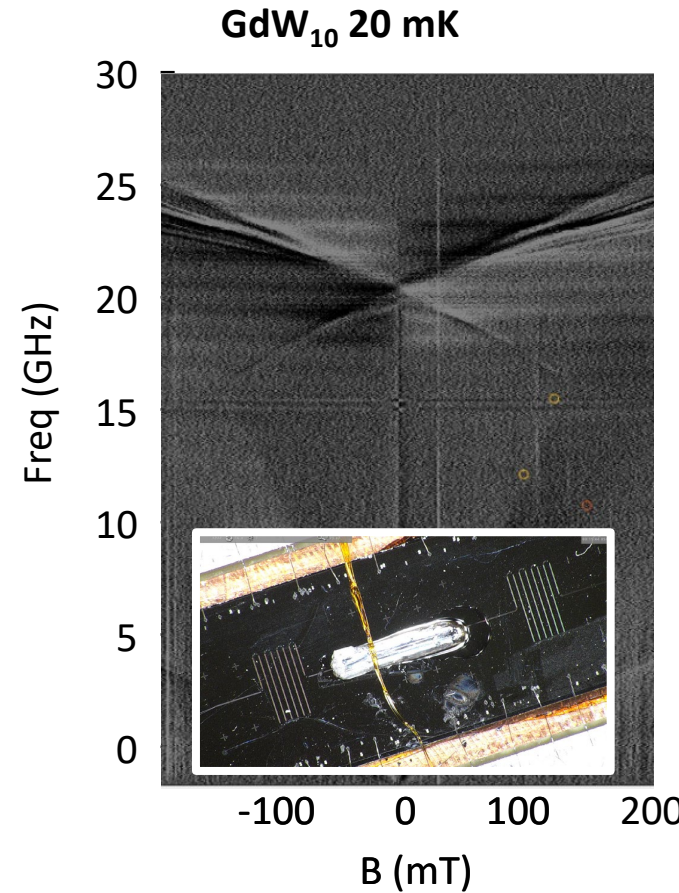
GdW10 a prototype spin qubit

GdW10 has a well defined easy axis. Hamiltonian: $H = B_2^0 \mathbf{O}_2^0 + B_4^4 \mathbf{O}_4^4 - g\mu_B \left[\frac{1}{2} (S_+ + S_-)(H_x + iH_y) + S_z H_z \right]$

$$B_2^0 = -0.059 \text{ K} \quad B_4^4 = 4 \times 10^{-4} \text{ K}$$



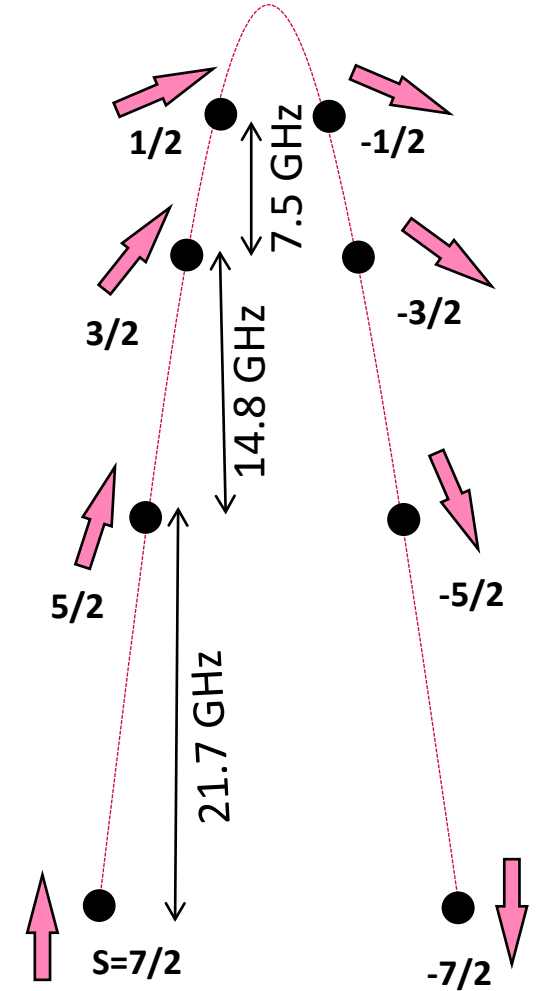
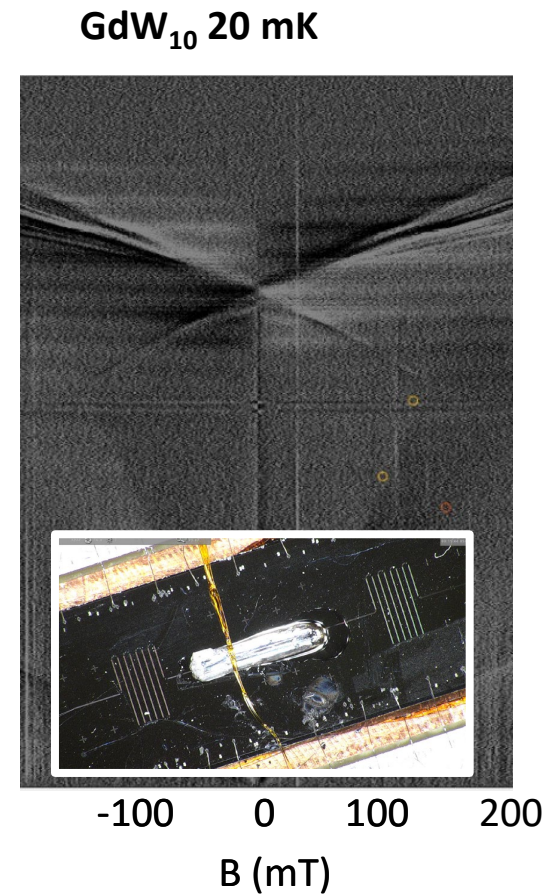
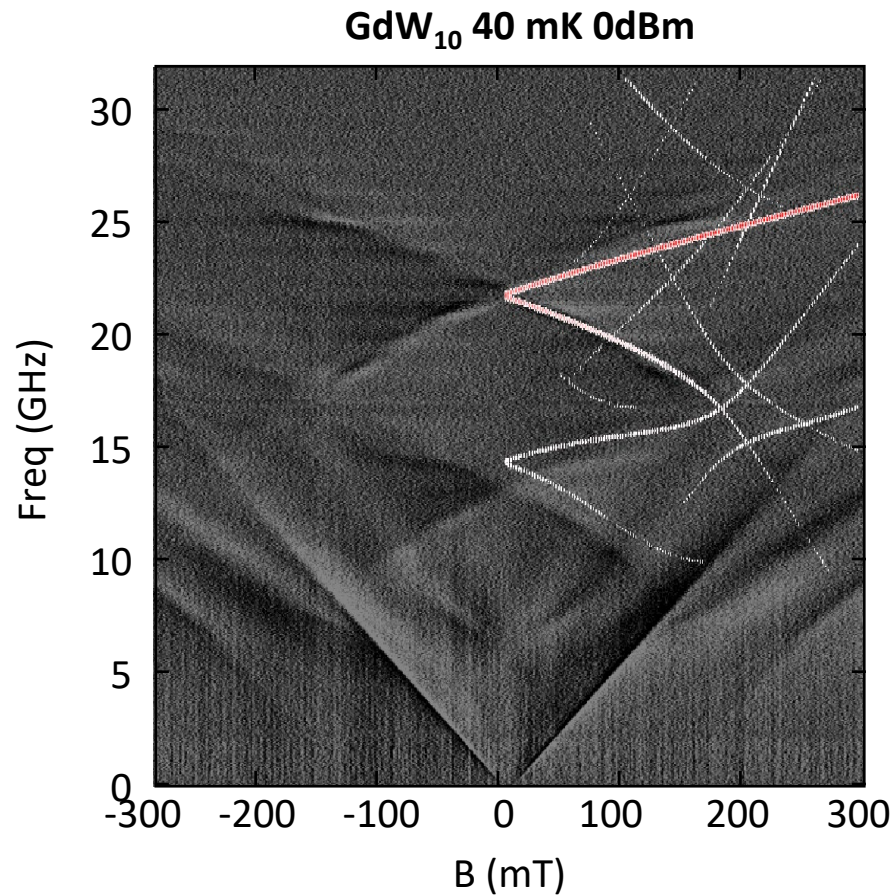
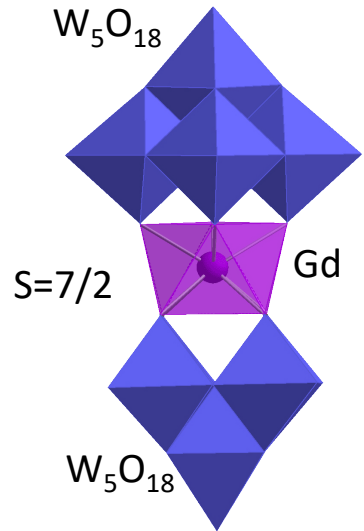
- Relatively large Zero Field Splitting of ~40 GHz (2 K)
- Dipolar order at 36 mK (very weak dipolar interactions)



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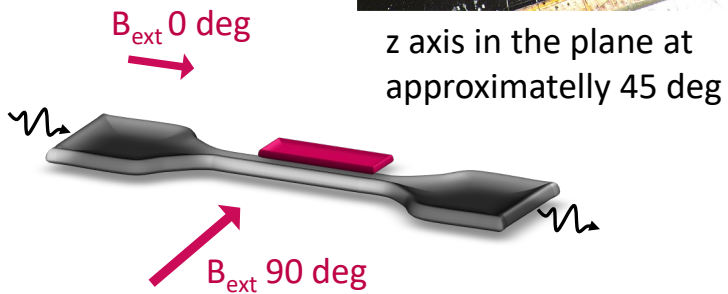
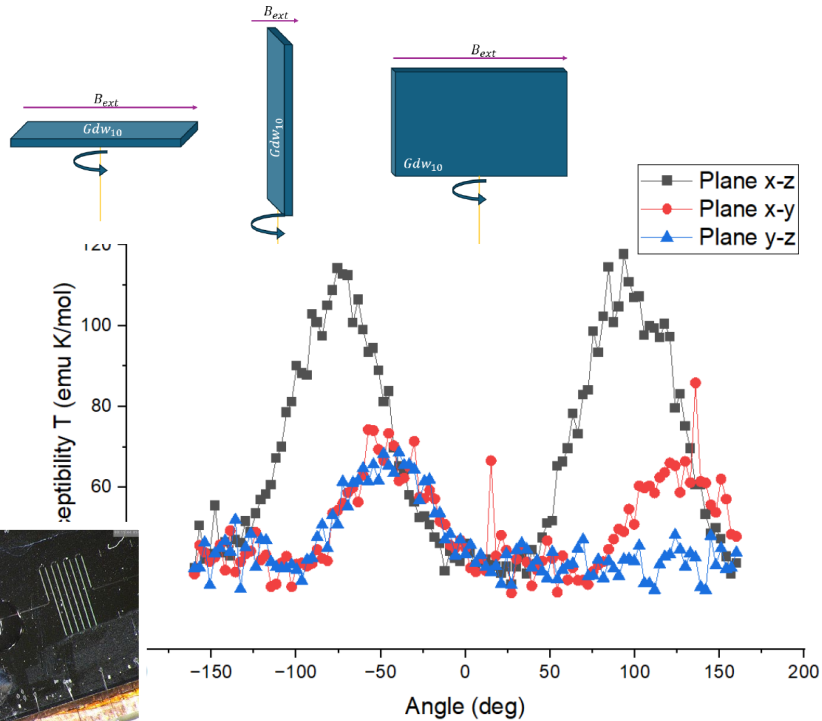
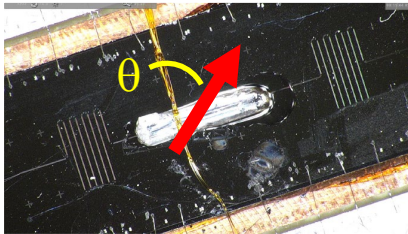
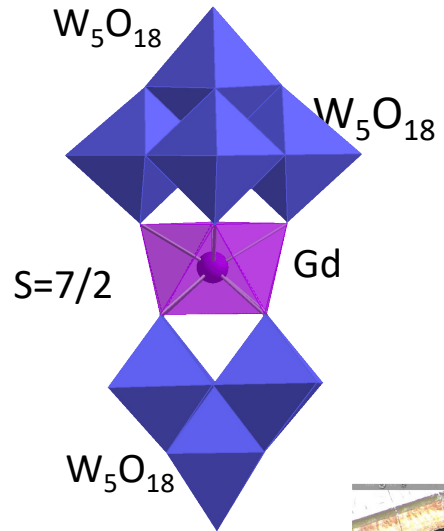
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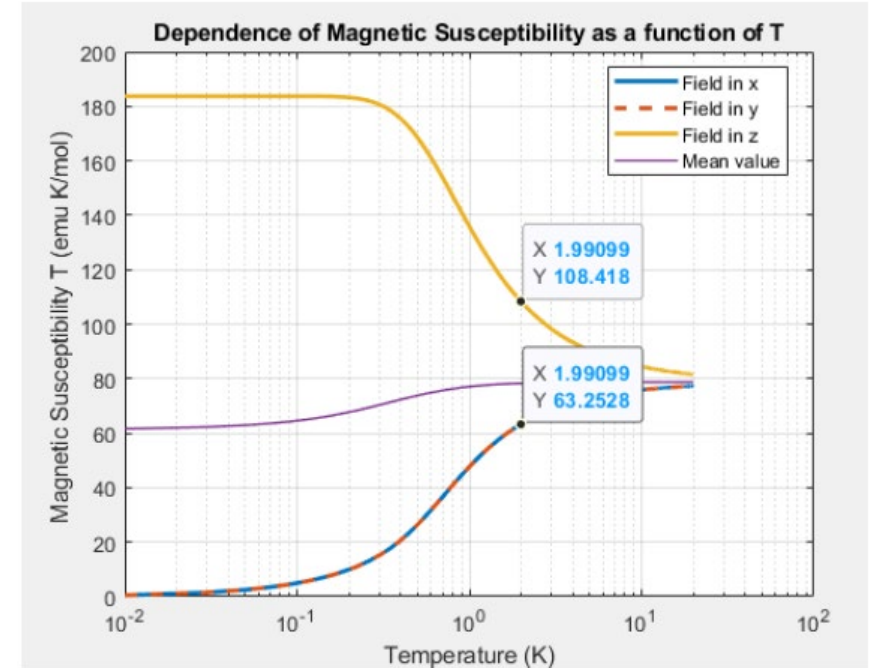
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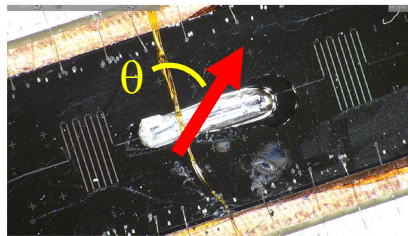
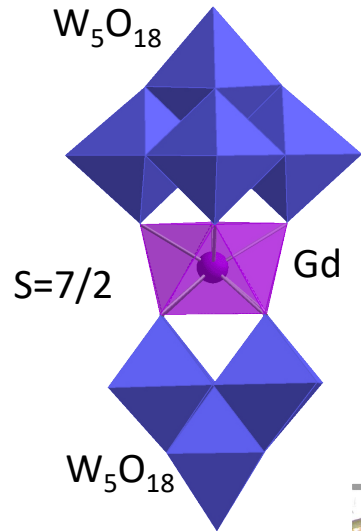
(collaboration with E. Coronado ICMOL)



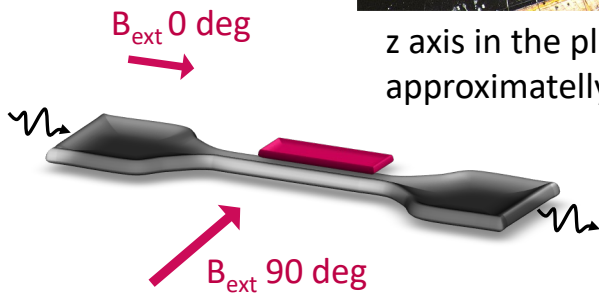
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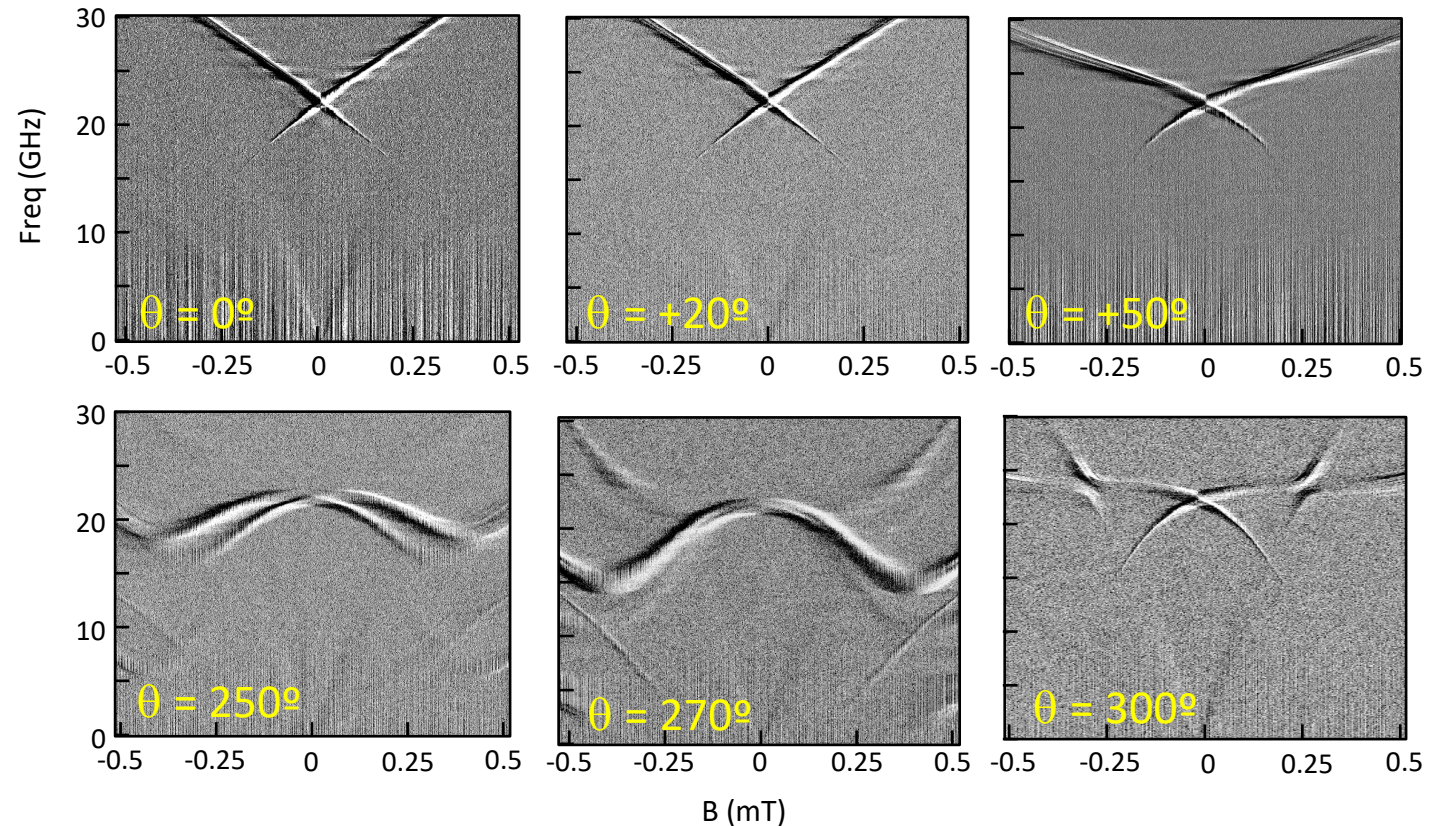
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z axis in the plane at approximately 45 deg



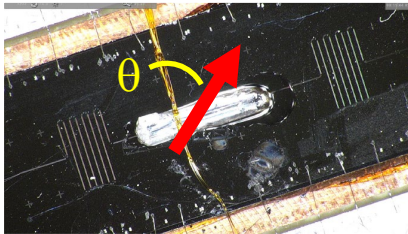
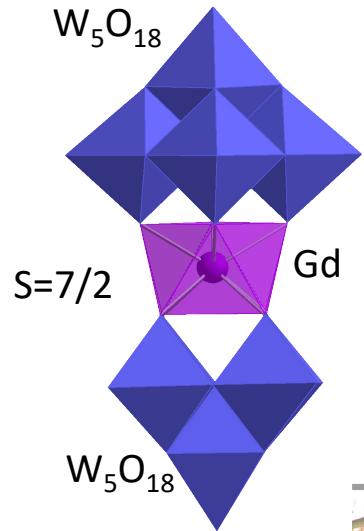
(collaboration with E. Coronado ICMOL)



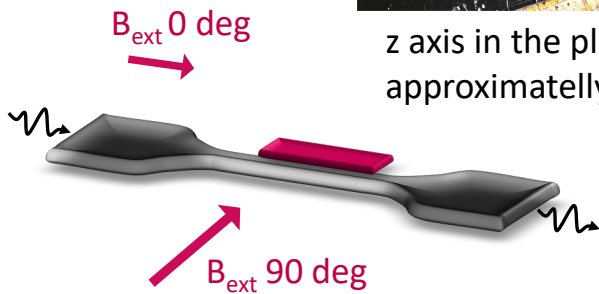
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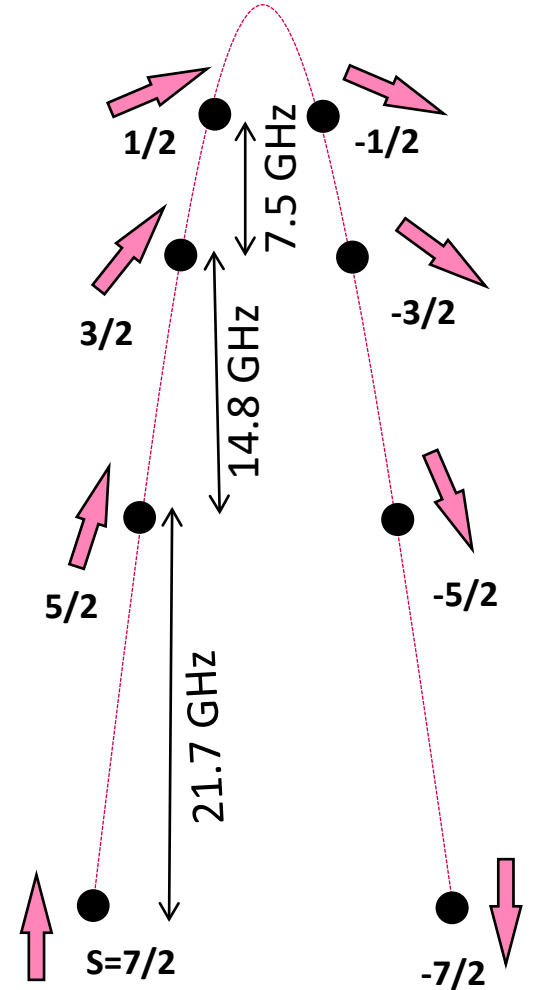
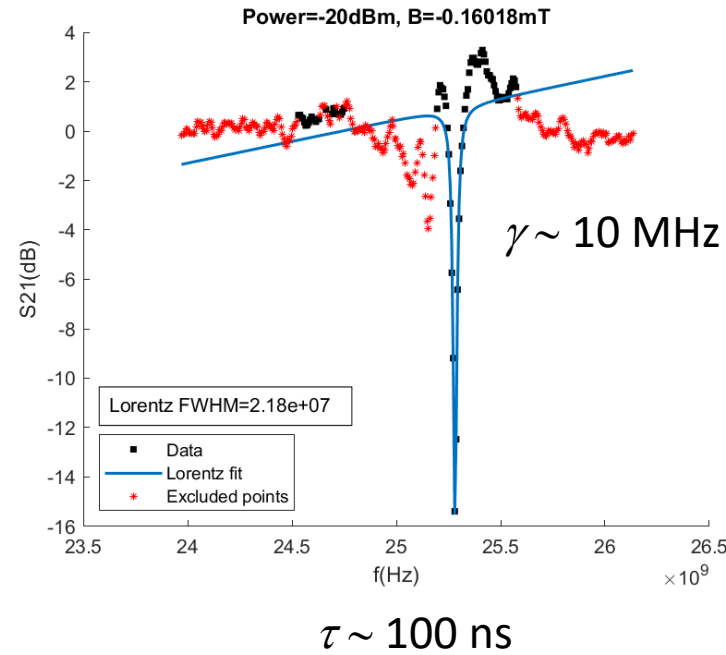
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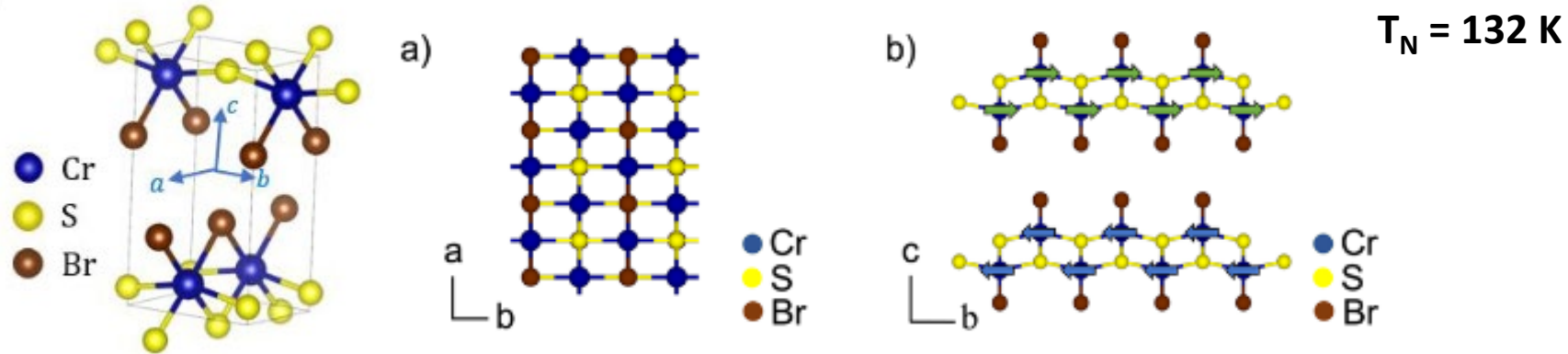


(collaboration with E. Coronado ICMOL)

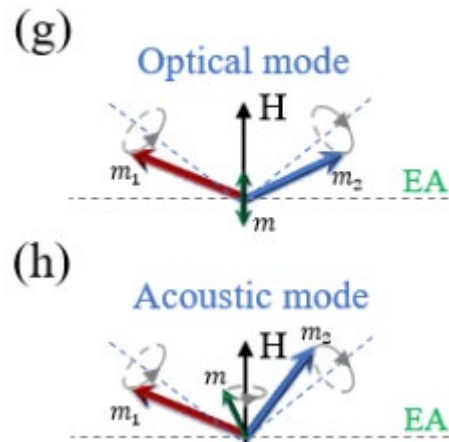


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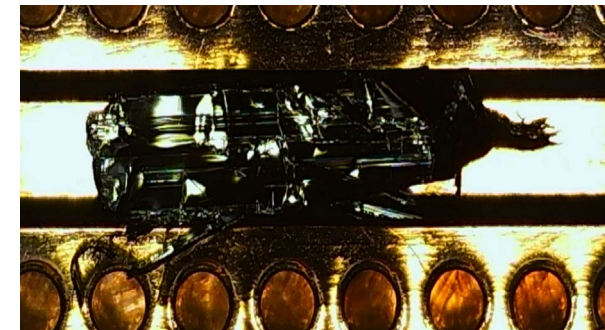
Van der Waals antiferromagnet: Chromium Sulfide Bromide



c axis is hard, the b axis is easy, a axis is intermediate



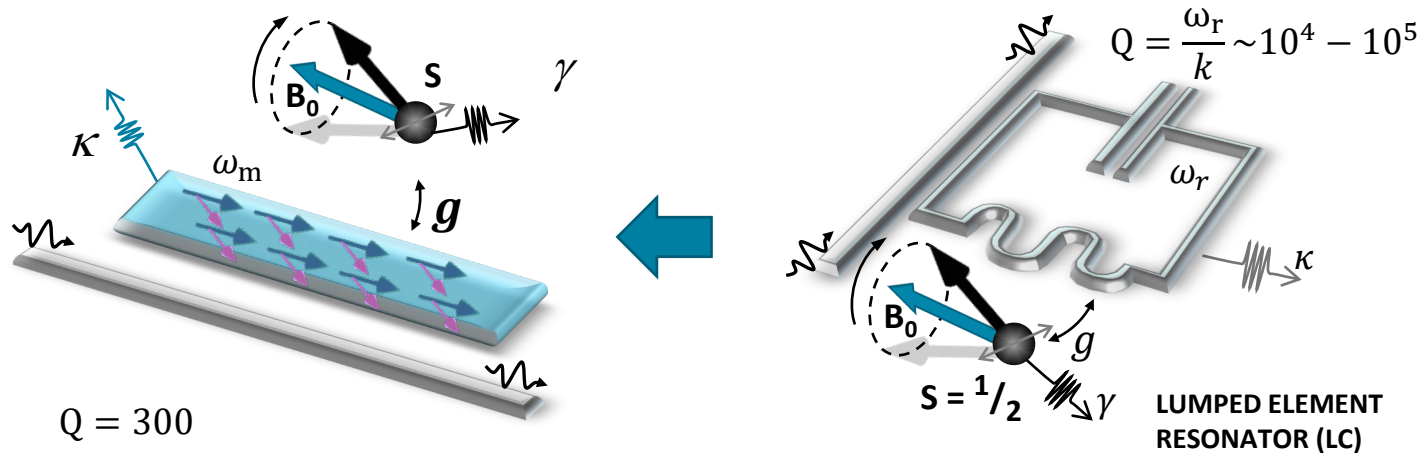
$B_{ex} = 0.395 \text{ T}$
 $B_{k,c} = 1.30 \text{ T}$ $B_{k,a} = 0.383 \text{ T}$



Two coupled LLG equations

$$\frac{d\hat{\mathbf{m}}_{1(2)}}{dt} = -\mu_0\gamma\hat{\mathbf{m}}_{1(2)} \times (\mathbf{H} - H_E\hat{\mathbf{m}}_{2(1)} - H_c(\hat{\mathbf{m}}_{1(2)} \cdot \hat{\mathbf{c}})\hat{\mathbf{c}} - H_a(\hat{\mathbf{m}}_{1(2)} \cdot \hat{\mathbf{a}})\hat{\mathbf{a}})$$

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Pérez-Bailón, (...) MJ M-P in preparation

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